



# At the limit: Quantum Computing

## Classical computer:

- the information is stored in classical bits, values 0,1
- usual operations NOT, AND, OR
- general purpose device



## Quantum computer:

- the information is stored in quantum bits (**qubits**)
- **unitary operations**, single- and two-qubit operations (XOR)
- powerful in calculating **specific tasks**

### Algorithms:

- prime factorization (Shor)
- data base search (Grover)
- random number generator
  - quantum simulation

**Quantum computation** is an interdisciplinary field, with contributions from mathematics, (theoretical) computer science, physics, and chemistry.

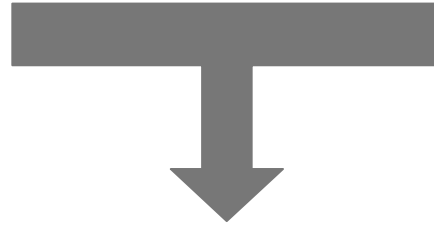
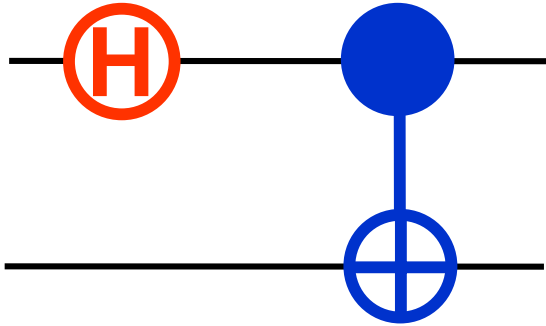
### Applications:

- quantum teleportation
- quantum cryptography

### Implementations:

- Quantum optics
  - **Solid state**
    - NMR

# Information meets Technology



**Information Science**

Recent developments in  
&

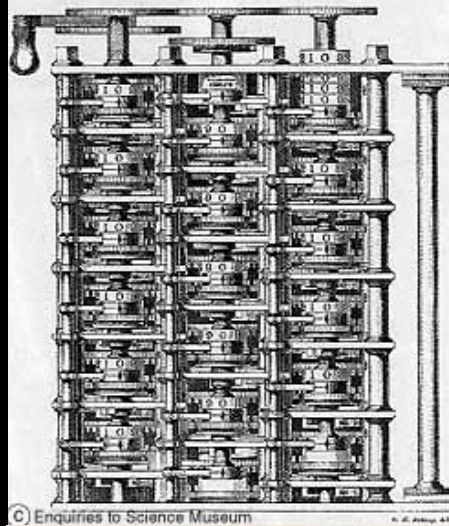
**Fundamental/  
Applied Physics**

merge to define a new common goal: the quantum computer

- **History:** from mechanics to nanoelectronics
- **Information Theory:** Turing machines & complexity
- **Quantum Mechanics:** superpositions & entanglement
- **Quantum Games:** no-cloning, cryptography, teleportation
- **Quantum Bits and Gates**
- **Quantum Algorithms:** Shor's period finder
- **Hardware:** superconducting phase qubits



Ada Byron,  
Lady Lovelace  
1815-1852

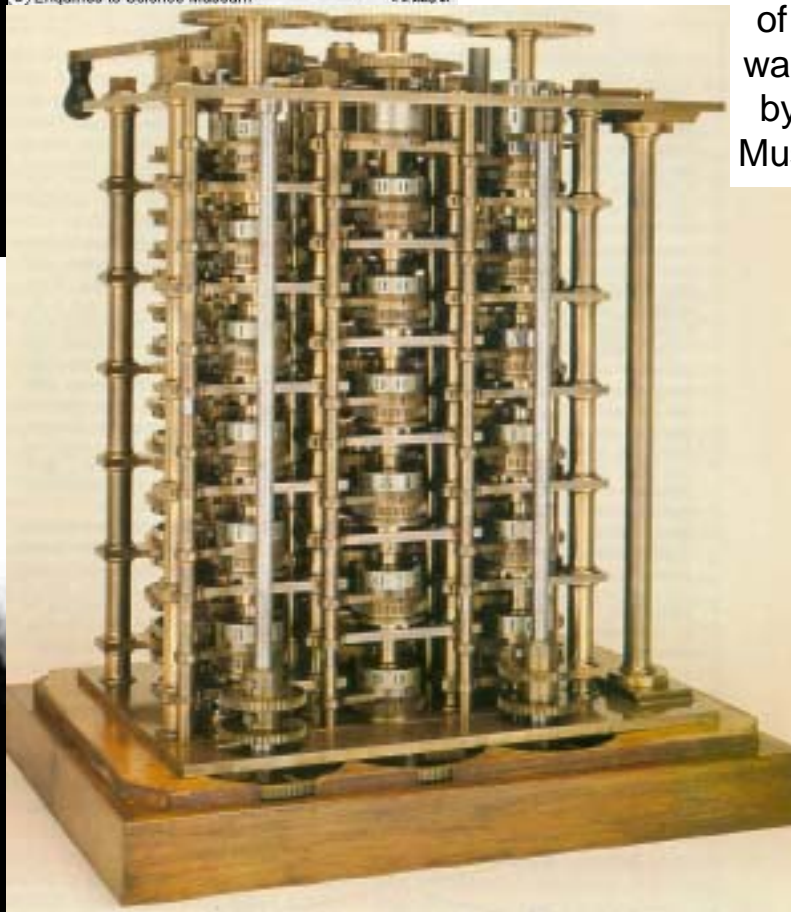


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# Mechanics

First 'Programmer'  
and  
Inventor of the  
Difference Engine  
1834

The 'full' version  
of this machine  
was built in 1991  
by the Science  
Museum, London



Enigma,  
cracked by  
Alan Turing  
with help of  
COLOSSUS



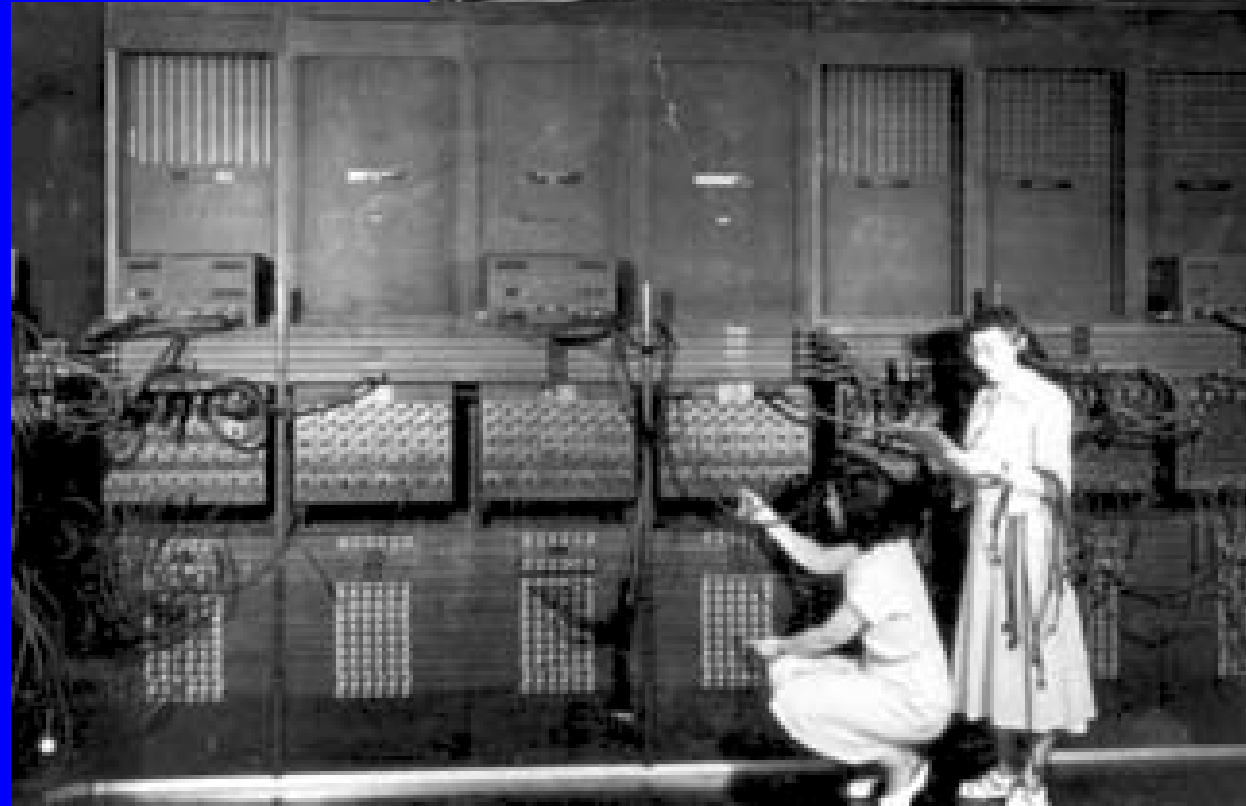
Charles Babbage  
1791-1871





# Electronics

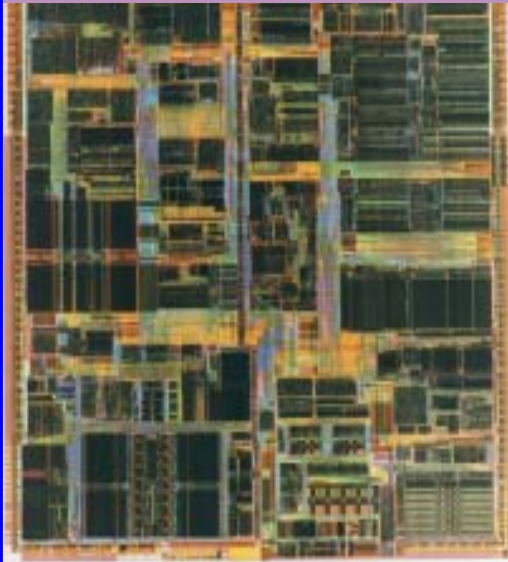
The ENIAC  
(Electronic Numerical  
Integrator and Computer)  
computer  
was built in 1946



Built at University of Pennsylvania,  
it included 18'000 tubes,  
weighed 30 tons,  
required 6 operators,  
and 160 m<sup>2</sup> of space.



Pentium Processor,  
1997, Intel



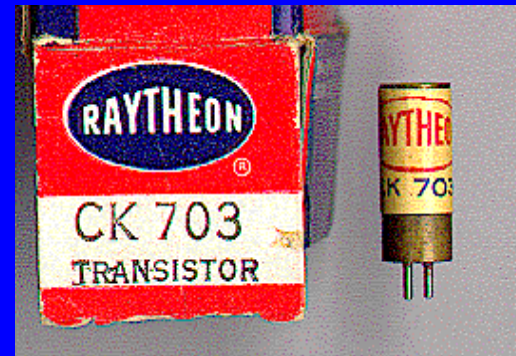
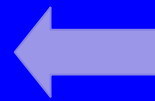
# Transistors



First Transistor, 1947  
Bell Laboratories  
Bardeen, Brattain,  
& Shockley



First Integrated Circuit, 1958  
Jack Kilby, Texas Instruments



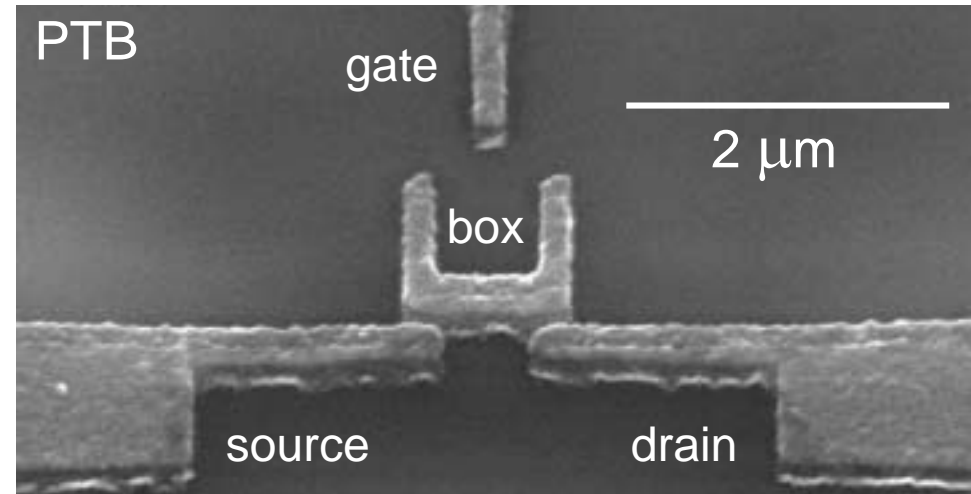
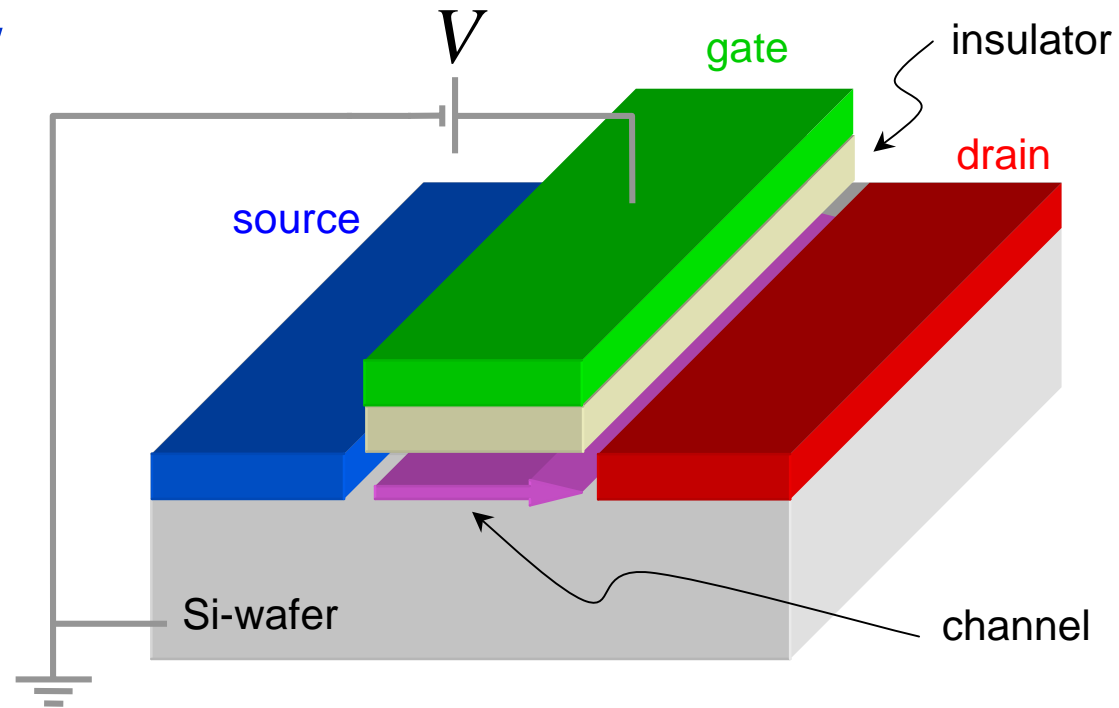
Packed Device



# Nanoscale Technology



Ultra-short channel Si-MOSFET,  
IBM  
0.5 μm wide, 0.1 μm channel



Single Electron Transistor (SET), Al-Technology

Switch a MOSFET with **1000** electrons,  
while a SET requires only **one!**

# Information Theory





Claude  
Shannon

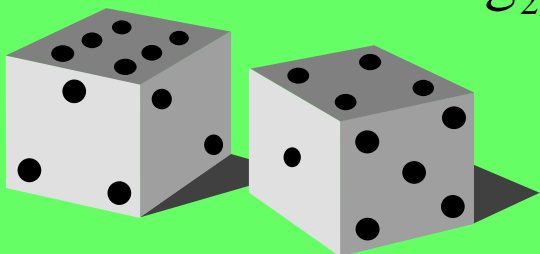
# Information Theory

Can be **quantified**: The random variable  $X$  distributed according to  $p(x)$  contains the information

$$S[p(x)] = -\sum_x p(x) \log_2 p(x)$$

E.g., in the process of throwing a dice one may gain the information

$$S = -\log_2 (1/6)$$



**Information**: is a **general concept**, similar to the concept of energy (appearing in many forms, e.g., mechanical, thermal, electrical,...).

Can be packed into **equivalent forms**:

0, 1      ↑, ↓      □, ■

This text is difficult

*Dieser Text ist schwierig*

Ce texte est difficile

**Information is physical** (Landauer, 1991):

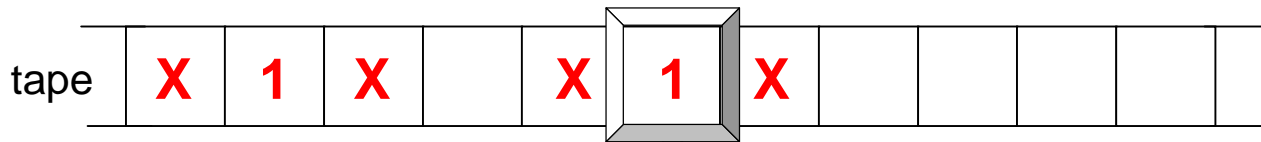
Need physical implementation to express and manipulate information; e.g., ink molecules on paper, charges in capacitors, currents in leads.

Rolf  
Landauer



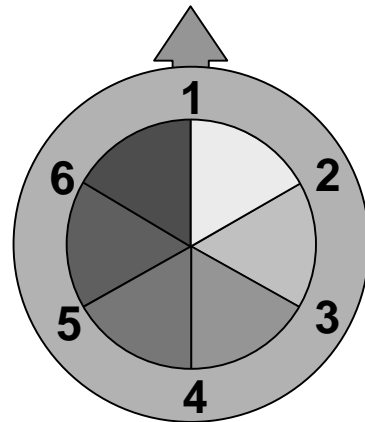
# Turing Machine (mid 1930)

machine



input

scanner



dial

program

		scanner		
		dial	X	1
Write Erase Move	1	W <sub>1</sub> 6	E 2	M <sub>R</sub> 1
	2	M <sub>R</sub> 2	E 3	Error
	3	M <sub>R</sub> 3	E 4	E 5
	4	M <sub>L</sub> 4	Error	M <sub>R</sub> 6
	5	M <sub>L</sub> 5	Error	M <sub>R</sub> 1
	6	W <sub>X</sub> 6	Done	M <sub>R</sub> 3



# A Universal Computer

**reproduces** the action of any other computer :

Let **T** be a Turing machine acting on an input **x**.  
There exists a universal machine **U** which takes **x**  
and a (binary) description **d [T]** as an input and  
reproduces the action of **T**,  $\mathbf{U}(\mathbf{d} [\mathbf{T}], \mathbf{x}) = \mathbf{T}(\mathbf{x})$ ,  
with polynomial effort in **d**.

Other models of computation,  
e.g., the **network model of computing**,  
are equivalent to the Turing model

concatenated logic gates  
acting on n-bit symbols

## Church-Turing Thesis (unproven)

Every function which would naturally be regarded  
as computable can be computed by a  
universal Turing machine.

This notion of universality allows us to  
classify computational problems

# Computational Complexity

An input  $\mathbf{x}$  is quantified via its information content  $\mathbf{L} = \log_2 \mathbf{x}$ .  
A calculation is characterized by the number  $\mathbf{s}$  of steps (logical gates) involved.

A problem is class **P** (efficiently solvable) if  $\mathbf{s}$  is polynomial in  $\mathbf{L}$ ,  $s \sim L^\mu$   
A problem is deemed 'hard' (not in P) if  $\mathbf{s}$  scales exponentially in  $\mathbf{L}$ ,  $s \sim \exp L^\nu$

A 'classic' hard problem is that of **prime factorization**:  
given a non-prime number  $\mathbf{N}$ , find its factors;  
the best known algorithm scales as  $\mathbf{s} \sim \exp(2 \mathbf{L}^{1/3} (\ln \mathbf{L})^{2/3})$ .

A modern computer can factor a 130-decimal-digits number ( $L = 300$ ) in a few weeks – days;



1827365426354265930284950398726453672819048374987653426354857645283905612849667483920396069782635471628694637109586756325221365901

doubling  $L$  would take **millions of years** to carry out this calculation.

**A quantum computer would do the job within minutes**

# Public Key Encryption

(Rivest, Shamir & Adleman, 1978)

$M$  = message

$s$ ,  $N = p \cdot q$ , (non-)public key

Encoding



$$E = M^s \pmod N$$

Decoding



$$M = E^{t(s,p,q)} \pmod N$$

A quantum computer would crack this encryption scheme



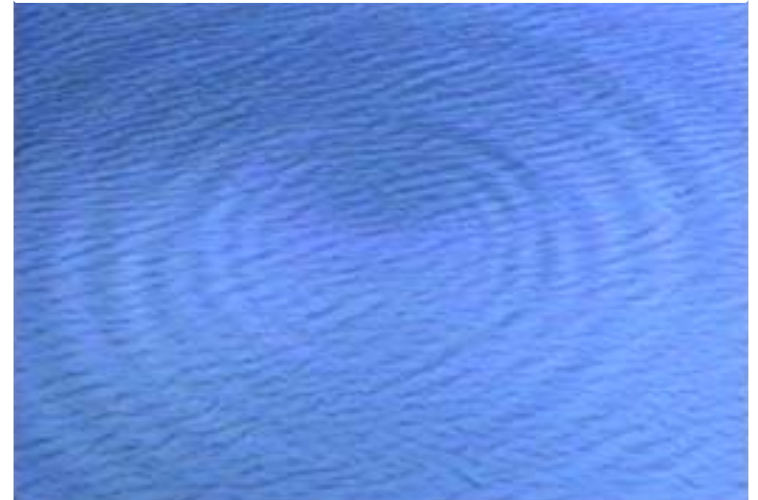
# Quantum Mechanics

# Quantum Mechanics

The following two elements of quantum mechanics are central to quantum computing

## **Superposition of states:**

A quantum degree of freedom is described through a wave function.

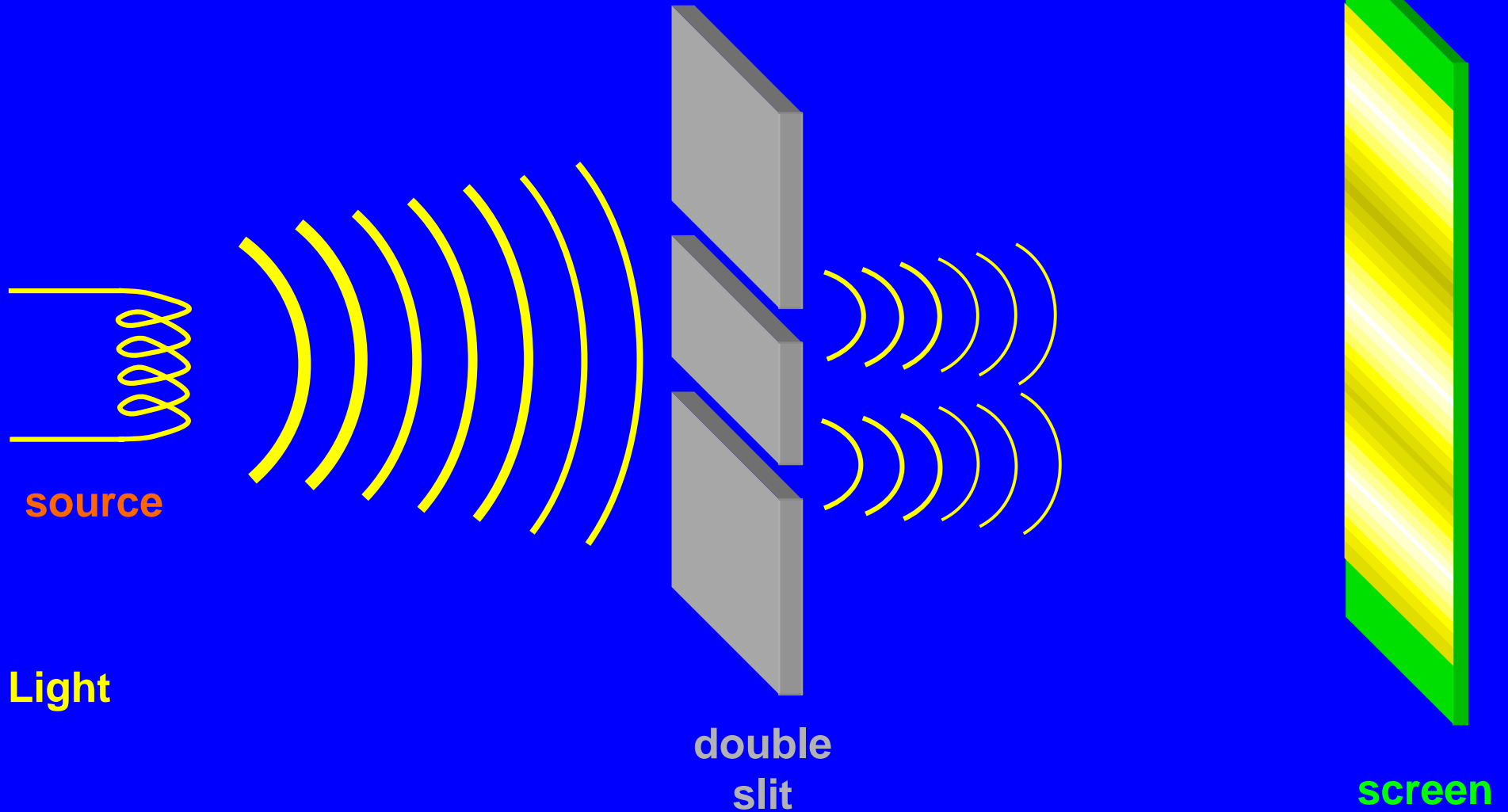


## **Entanglement of states:**

Two quantum degrees of freedom can exhibit stronger correlations than any classical system.

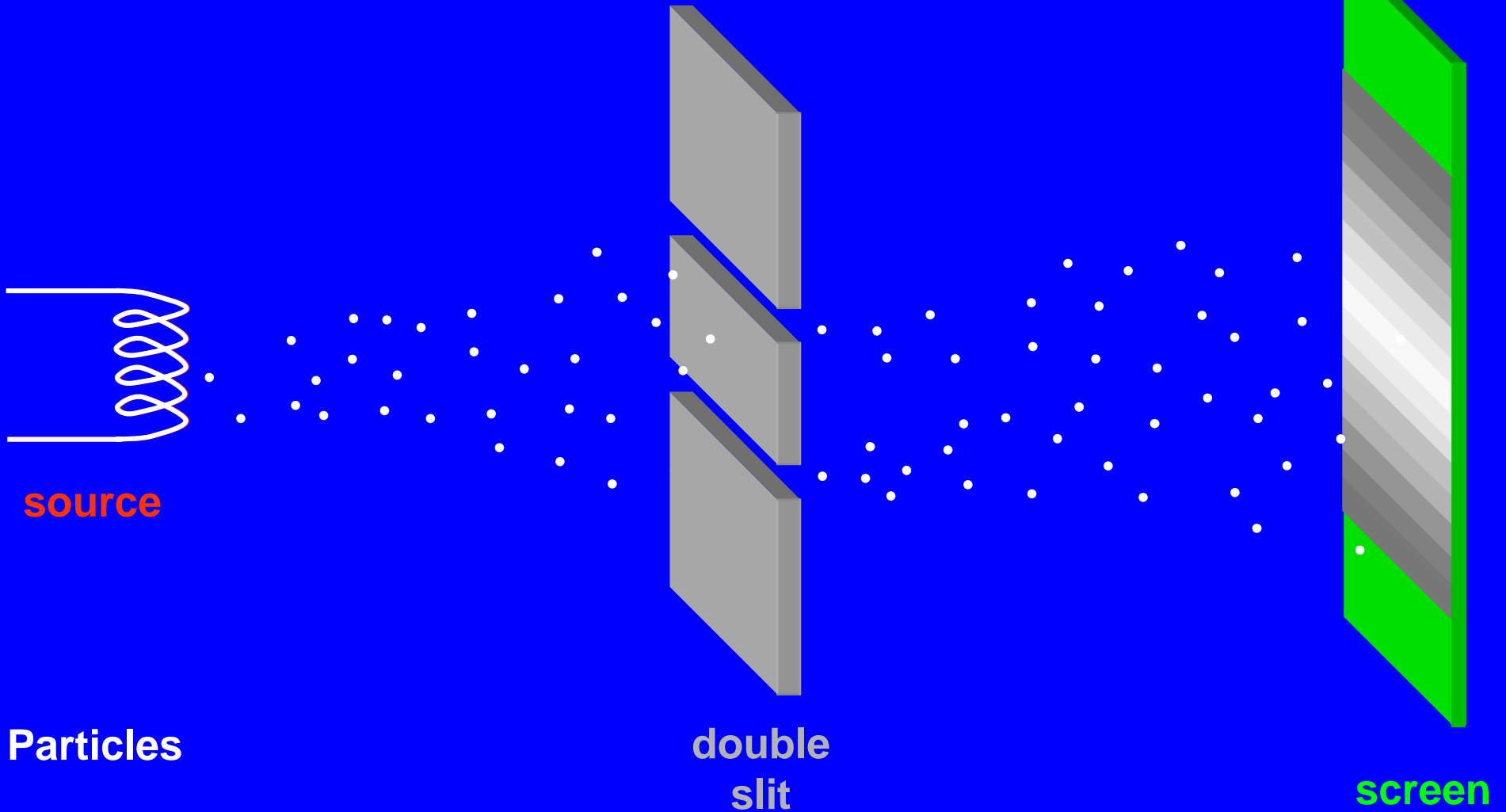
# Waves

## The double slit experiment

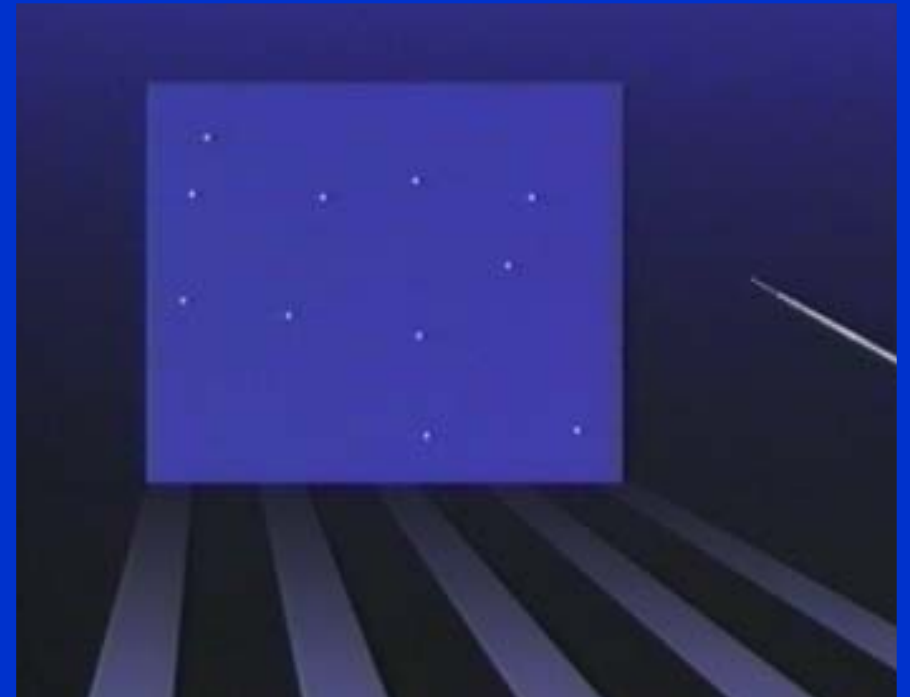


# Waves

## The double slit experiment

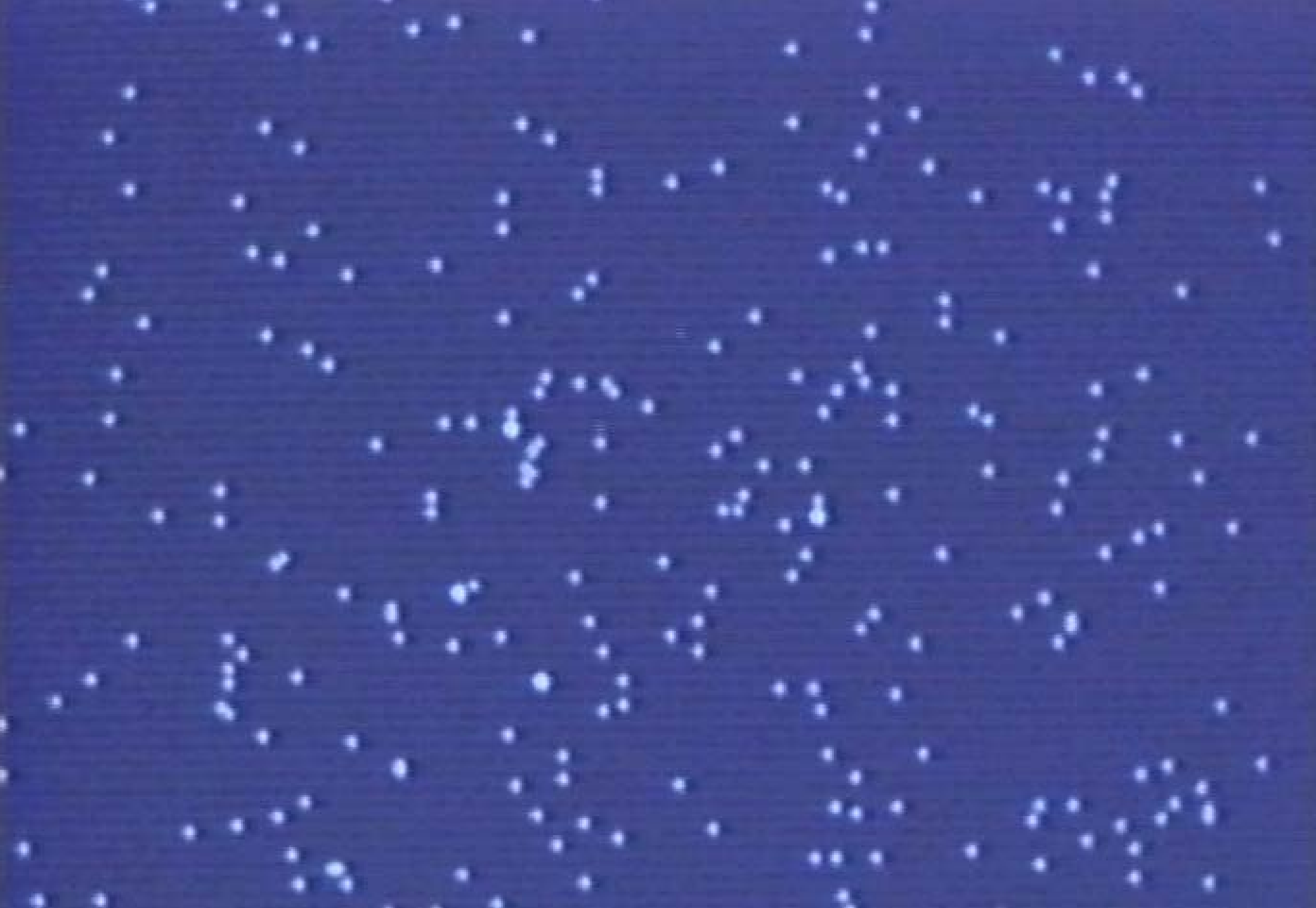


# Particle-Wave Duality



- $V = 50 \text{ kV}$  electrons
- $\lambda = 0.05 \text{ \AA}$
- $10^3$  electron / sec
- source–screen distance = 1.5 m
- average electron distance 150 km
- size of electronic wave packet  $1 \text{ }\mu\text{m}$
- total exposure time 20 min



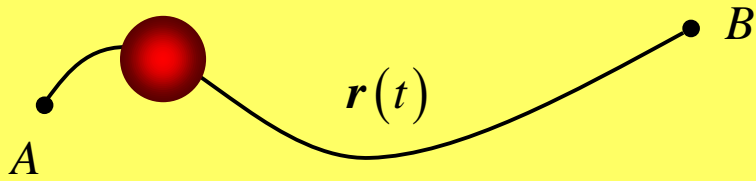


# Classical

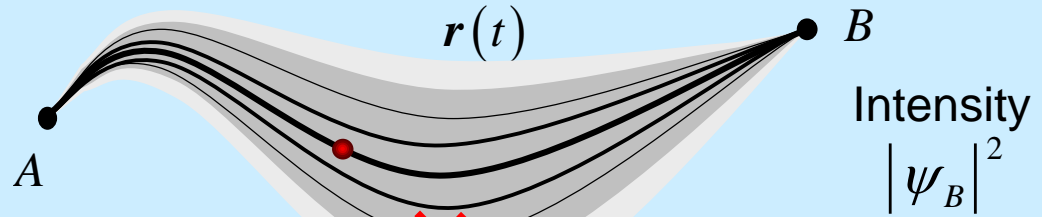
vs

# quantum mechanics

A quantum particle probes all trajectories  $r(t)$  from its initial starting point A to its final end point B.

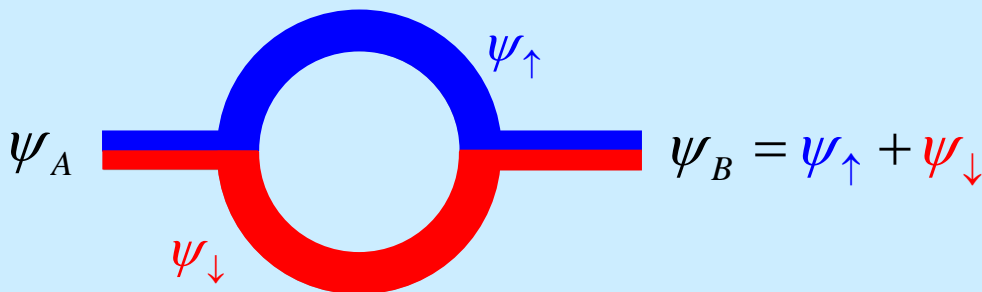


A classical particle follows its trajectory  $r(t)$  from its initial starting point A to its final end point B.



Intensity  $|\psi_B|^2$   
 quantum parallelism is fragile: external perturbations easily destroy the quantum coherent state.

If we offer a particle only two paths then the wave function is the **superposition** of two amplitudes:



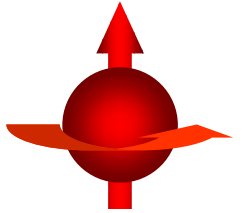
Its wave function is a sum over amplitudes over all paths,

$$\psi_B = \sum_{\text{paths}} \exp[i\varphi_{\text{path}}] \psi_A$$

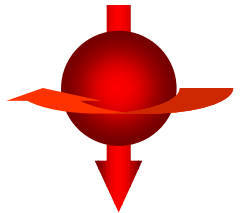
phase of the path,  $\varphi_{\text{path}} = S_A^B(\text{path}),$   
 $S_A^B = \frac{1}{\hbar} \int_{r_A,0}^{r_B,t} dt \left[ m \frac{\dot{r}^2}{2} - V(r) \right].$

# Spins

Many elementary particles (such as electrons) carry a spin, an internal angular momentum taking on two values:

$$\psi_{\uparrow} = |\uparrow\rangle$$


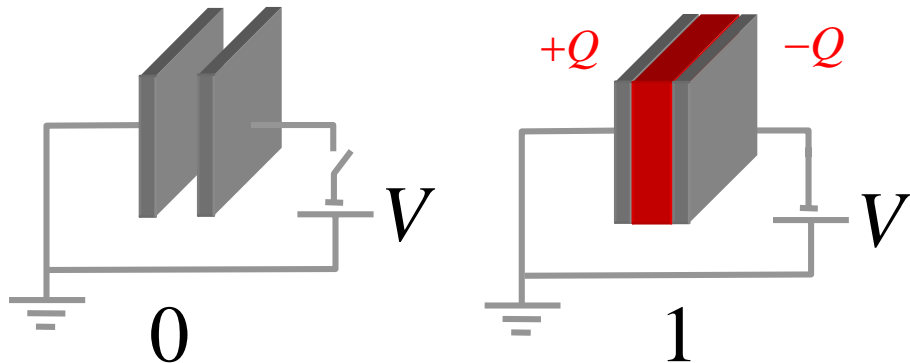
A spin system is a generic two-level system which is the generalization of a **classical bit** to a **quantum bit** or **qubit**:

$$\psi_{\downarrow} = |\downarrow\rangle$$


Classical bit

0, 1

Physical realization via a charged/uncharged capacitor

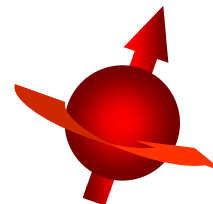


Quantum bit

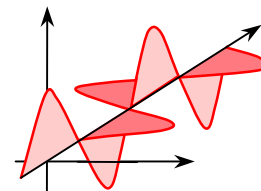
$$|a, \varphi\rangle = \left[ |0\rangle + a e^{i\varphi} |1\rangle \right] / \sqrt{1 + a^2}$$

Physical realization via a two-level system

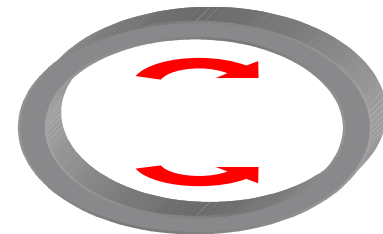
spin



polarization



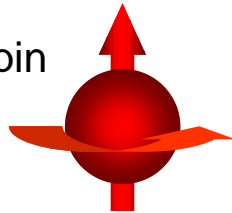
ring-current



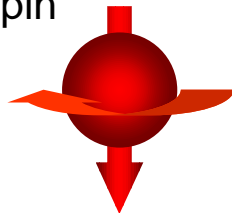
# Quantization axis: preparing spins

Choose an axis,  
e.g., the  $z$ -axis,  
and measure the spin,  
we will find either

an up-spin



or a down-spin



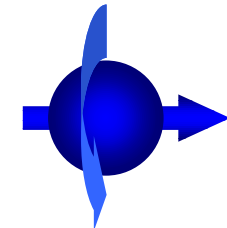
quantization  
axis

The reason is, that a spin prepared along a definite (quantization) axis, is a (50/50) superposition of states when viewed from another axis.

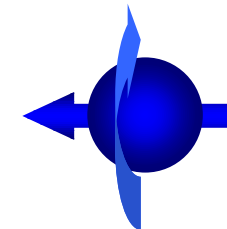
Let us assume we measured an up spin.  
We can re-measure the spin to check if it is still up and it will always be

But then, let us measure the spin along another axis, say the  $x$ -axis; we will find a new result, namely, half of the time the spin will point along

the positive direction



the negative direction



$$|\uparrow_z\rangle = [|\rightarrow_x\rangle + |\leftarrow_x\rangle] / \sqrt{2}$$

# Cryptography

We can make use of spin / two-level systems for cryptography: Alice and Bob wish to exchange a secret key, a sequence of 0 and 1. In order to do so, Alice sends Bob a sequence of spins which she polarized either along the  $z$ - or  $x$ - axis.

$$|\uparrow\rangle = \begin{array}{c} \uparrow \\ \text{red sphere} \\ \downarrow \end{array} \quad \begin{array}{c} z \\ \uparrow \end{array} \quad \begin{array}{c} \downarrow \\ \text{red sphere} \\ \uparrow \end{array} = |\downarrow\rangle$$

$$|+\rangle = \begin{array}{c} \text{blue sphere} \\ \rightarrow \end{array} \quad \begin{array}{c} \leftarrow \\ \text{blue sphere} \\ \rightarrow \end{array}$$

$$[|\uparrow\rangle + |\downarrow\rangle] / \sqrt{2} = \begin{array}{c} \text{blue sphere} \\ \rightarrow \end{array} \quad \begin{array}{c} \leftarrow \\ \text{blue sphere} \\ \rightarrow \end{array} \quad \begin{array}{c} x \\ \rightarrow \end{array}$$

$$|-\rangle = [|\uparrow\rangle - |\downarrow\rangle] / \sqrt{2}$$

$$\uparrow + - + \downarrow \downarrow - = \begin{array}{c} \uparrow \\ \text{red sphere} \\ \downarrow \end{array} \quad \begin{array}{c} \rightarrow \\ \text{blue sphere} \\ \leftarrow \end{array} \quad \begin{array}{c} \leftarrow \\ \text{blue sphere} \\ \rightarrow \end{array} \quad \begin{array}{c} \rightarrow \\ \text{blue sphere} \\ \leftarrow \end{array} \quad \begin{array}{c} \downarrow \\ \text{red sphere} \\ \uparrow \end{array} \quad \begin{array}{c} \downarrow \\ \text{red sphere} \\ \uparrow \end{array} \quad \begin{array}{c} \rightarrow \\ \text{blue sphere} \\ \leftarrow \end{array}$$

Bob measures the spins he receives from Alice, choosing randomly among the  $x$ - and  $z$ - axis

$$\begin{array}{c} z \\ \uparrow \end{array} \quad \begin{array}{c} \rightarrow \\ x \end{array} \quad \begin{array}{c} z \\ \uparrow \end{array} \quad \begin{array}{c} \rightarrow \\ x \end{array} \quad \begin{array}{c} \rightarrow \\ x \end{array} \quad \begin{array}{c} z \\ \uparrow \end{array} \quad \begin{array}{c} z \\ \uparrow \end{array}$$

$$\begin{array}{c} \uparrow \\ \text{red sphere} \\ \downarrow \end{array} \quad \begin{array}{c} \rightarrow \\ \text{blue sphere} \\ \leftarrow \end{array} \quad \begin{array}{c} \uparrow \\ \text{green sphere} \\ \downarrow \end{array} \quad \begin{array}{c} \rightarrow \\ \text{blue sphere} \\ \leftarrow \end{array} \quad \begin{array}{c} \leftarrow \\ \text{green sphere} \\ \rightarrow \end{array} \quad \begin{array}{c} \downarrow \\ \text{red sphere} \\ \uparrow \end{array} \quad \begin{array}{c} \downarrow \\ \text{green sphere} \\ \uparrow \end{array} = \uparrow + \uparrow + - \downarrow \downarrow$$

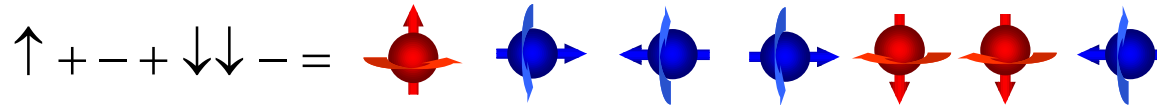
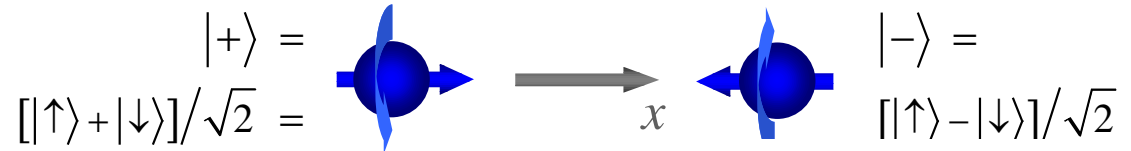
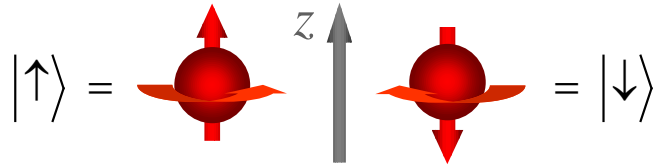


When Alice publicly announces along which axes she has prepared the individual spins, Bob can identify the **bad** spins and inform Alice – in the end they possess a ‘good’ key made from  $|\uparrow\rangle, |+\rangle \Rightarrow 0$  and  $|\downarrow\rangle, |-\rangle \Rightarrow 1$ .

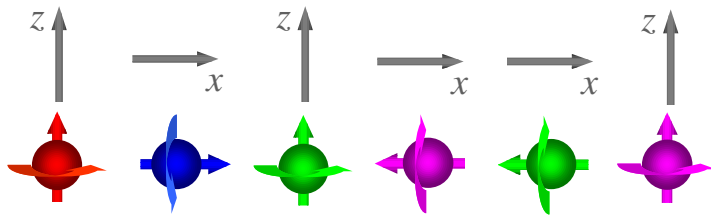


# Cryptography

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Bob measures the spin he receives from Alice, choosing randomly among the  $x$ - and  $z$ - axis



And when Eve interferes to learn about Alice's and Bob's secret key, she will spoil the sequence such that Bob and Alice can detect her presence.

Alice and Bob share a probabilistically secure key 

# Classical & quantum gates I

The possibilities to manipulate a classical bit are quite limited:  
The NOT-gate simply interchanges the two values 0 and 1 of the classical bit.

<b>i</b>	<b>f</b>
<b><i>0</i></b>	<b><i>1</i></b>
<b><i>1</i></b>	<b><i>0</i></b>

On the other hand, manipulation of a quantum bit is much richer!

# Classical & quantum gates I

A single qubit is manipulated via **unitary transformations**  $U(t)$ , which is just the usual **time evolution of quantum mechanics**:

Schrödinger equation  $\rightarrow$   $i\hbar \partial_t \psi(t) = H \psi(t), \quad H = H^\dagger,$   $\leftarrow$  Hamiltonian

$$\psi(t) = \exp[-iH t/\hbar] \psi(0)$$

$U(t)$   $\rightarrow$

For a spin / two-level system we can perform rotations around the  $x$  -,  $y$  -, and  $z$  - axis; placing the spin  $S$  (with magnetic moment  $\mu$ ) into a magnetic field  $H$ , the Hamiltonian

$$H = -\mu S \cdot H$$

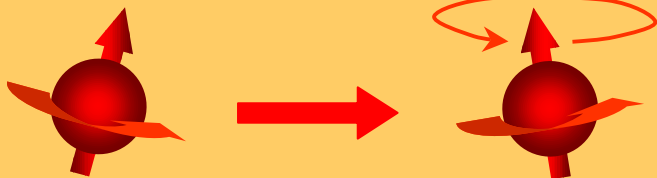
produces the desired rotation. E.g., with

$H = H_z$  we obtain the time evolution

$$U = \begin{pmatrix} e^{-i\mu H_z t/2} & 0 \\ 0 & e^{i\mu H_z t/2} \end{pmatrix},$$

phase shifter

$$|\psi(t)\rangle = [e^{-i\mu H_z t/2} |\uparrow\rangle + a e^{i\mu H_z t/2} |\downarrow\rangle] / \sqrt{1+a^2}.$$



$H = H_x$  we obtain the time evolution

$$U = \begin{pmatrix} \cos \mu H_x t & i \sin \mu H_x t \\ i \sin \mu H_x t & \cos \mu H_x t \end{pmatrix},$$

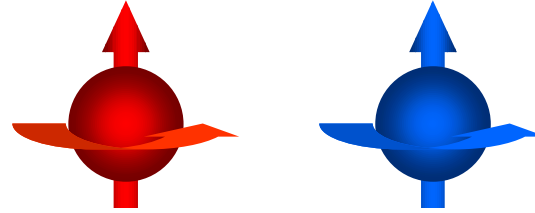
amplitude shifter ( $a = 0$ )

$$|\psi(t)\rangle = \cos \mu H_x t [|\uparrow\rangle + i \tan \mu H_x t |\downarrow\rangle].$$



# Entanglement

Consider two spins:



They can appear in a superposition of four states

$$|\uparrow\uparrow\rangle, |\uparrow\downarrow\rangle, |\downarrow\uparrow\rangle, |\downarrow\downarrow\rangle$$

The singlet state

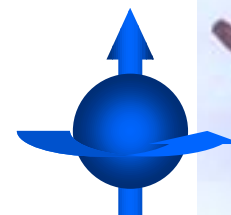
is an entangled state with astonishing properties:

$$|00\rangle = [|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle] / \sqrt{2}$$

After preparation of the singlet state (through having the two spins interact with one another) let us separate the two spins and give Alice and Bob each one of the two spins:

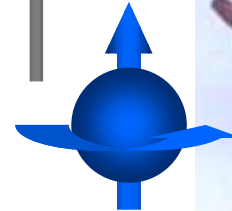
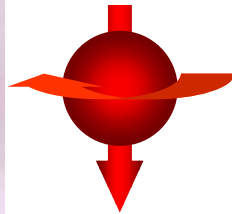


Alice and Bob now can experiment with their spins, e.g., measure their direction along the  $x$ -,  $y$ -, or  $z$ - axis, and they will find a few astonishing results!



# Einstein-Podolsky-Rosen

Have Alice measure her spin along an axis, say the  $z$ -axis. She has a fifty-fifty chance to measure a spin up – assume  $z$  she really finds her spin pointing up.



Next, have Bob measure his spin; although he may be arbitrarily far away from Alice, he will find his spin point down.

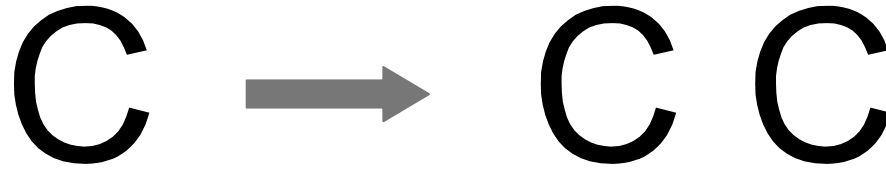
This is the **EPR-Paradox:**  
**quantum mechanics is non-local.**

The repetition of this experiment always gives the same result – the two entangled spins are fully correlated. A more careful analysis shows, that whenever Alice and Bob measure their spins along the directions  $\theta_A$  and  $\theta_B$  they will find them correlated to a degree  $\sin^2[(\theta_A - \theta_B)/2]$  – there is no classical process which will deliver such a high degree of correlations – here is a process which a classical computer cannot simulate!



# No-Cloning

The ability to copy a classical bit is widely exploited in computers and algorithms.



**There are no quantum copying machines: a qubit cannot be copied (cloned).**

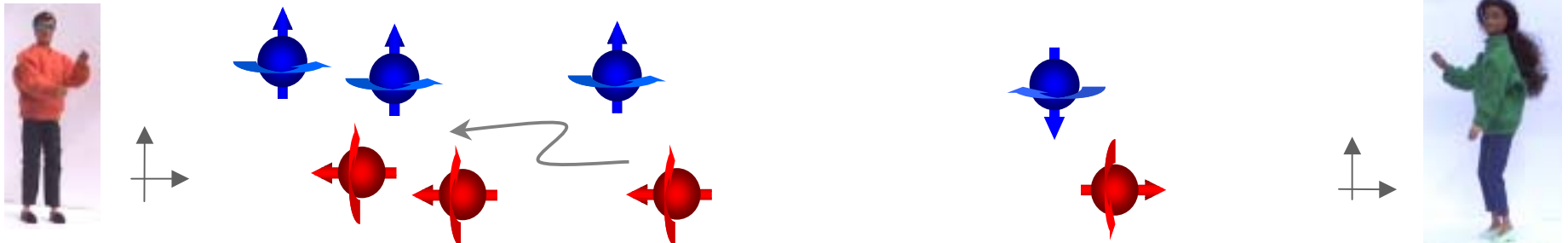
Assume the contrary, then there exists a unitary operator  $U$  independent of  $|\alpha\rangle$  and  $|\beta\rangle$  which produces copies of  $|\alpha\rangle$  and  $|\beta\rangle$ ,

$$U|\alpha 0\rangle = |\alpha\alpha\rangle \quad \text{and} \quad U|\beta 0\rangle = |\beta\beta\rangle.$$

But when we try to copy  $|\gamma\rangle = [|\alpha\rangle + |\beta\rangle]/\sqrt{2}$  our quantum copying machine  $U$  fails,

$$U|\gamma 0\rangle = [|\alpha\alpha\rangle + |\beta\beta\rangle]/\sqrt{2} \neq |\gamma\gamma\rangle.$$

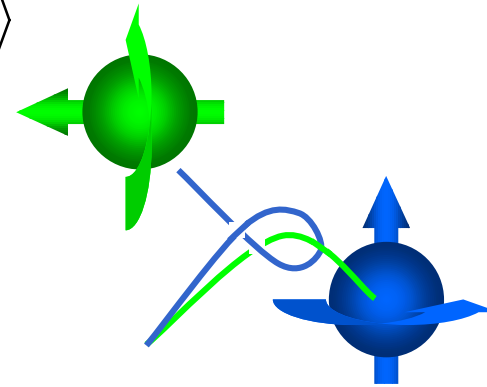
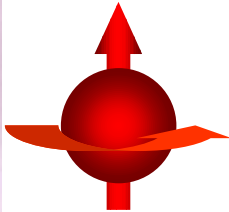
The combination of the EPR Paradox and the No-Cloning Theorem rescues the consistency between quantum mechanics and special relativity: For, if Bob could draw copies of entangled spins then Alice could communicate with Bob via a faster-than-light channel.



# Teleportation

Though a qubit  $|\phi\rangle = a|\uparrow\rangle + b|\downarrow\rangle$  cannot be copied, it can be teleported, i.e., it can be made to vanish at one place and reappear at another place. In order to do so, Alice and Bob share an entangled state  $[|\uparrow\uparrow\rangle + |\downarrow\downarrow\rangle]/\sqrt{2}$ .

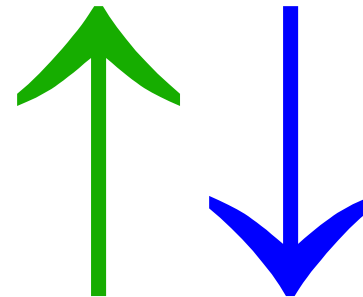
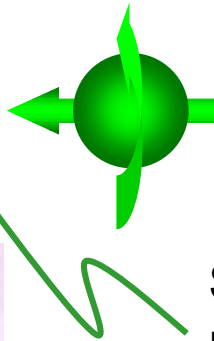
Next, Alice entangles her spin  $\uparrow$  with the unknown state  $|\phi\rangle$  and then measures what state her two spins are in.



# Teleportation

Though a qubit  $|\phi\rangle = a|\uparrow\rangle + b|\downarrow\rangle$  cannot be copied, it can be teleported, i.e., it can be made to vanish at one place and reappear at another place. In order to do so, Alice and Bob share an entangled state  $[|\uparrow\uparrow\rangle + |\downarrow\downarrow\rangle]/\sqrt{2}$ .

Next, Alice entangles her spin  $\uparrow$  with the unknown state  $|\phi\rangle$  and then measures what state her two spins are in.



She tells Bob, which of the four possible results  $\uparrow\uparrow, \uparrow\downarrow, \downarrow\uparrow, \downarrow\downarrow$  she has found and Bob carries out the appropriate rotation of his spin  $\uparrow$  by  $\pi$  around the  $I$ ,  $x$ -,  $z$ -, or  $y$ -axis.

As a result, Bob ends up with his spin in the state  $|\phi\rangle = a|\uparrow\rangle + b|\downarrow\rangle$ .



# Classical & quantum gates II

The combination of the classical gates allows us to construct all manipulations on classical bits.

NOT

i	f
0	1
1	0

AND

i	i	f
0	0	0
0	1	0
1	0	0
1	1	1

OR

i	i	f
0	0	0
0	1	1
1	0	1
1	1	1

irreversible

Is there a set of universal quantum gates ?  
How does such a set look like ?


## Single-qubit gates: Rotations

$$U = \begin{pmatrix} \cos(\theta/2) & -ie^{-i\phi} \sin(\theta/2) \\ -ie^{i\phi} \sin(\theta/2) & \cos(\theta/2) \end{pmatrix},$$

amplitude shifter

phase shifter

$$|\psi\rangle = \frac{1}{\sqrt{1+a^2}} \left[ |0\rangle + a e^{i\phi} |1\rangle \right]$$

Hadamard (basis change): 

$$H = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}; \quad \begin{matrix} |0\rangle \\ |1\rangle \end{matrix} \Leftrightarrow \begin{cases} |+\rangle = [|0\rangle + |1\rangle]/\sqrt{2}, \\ |-\rangle = [|0\rangle - |1\rangle]/\sqrt{2}. \end{cases}$$

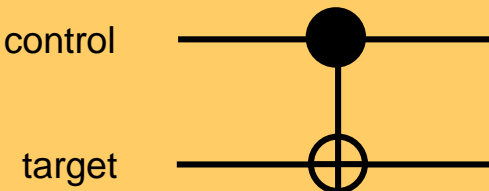
## Two-qubit gate: XOR (CNOT)

The target flips if the control is on 1

$$U = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{pmatrix}$$

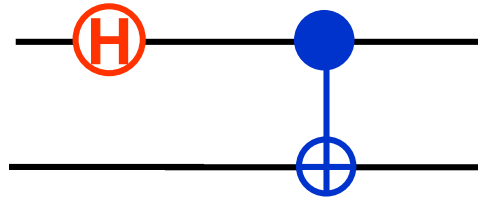
i	i	f
0	0	0
0	1	1
1	0	1
1	1	0

control



target

# Entangling two qubits



$$\mathbf{H}|00\rangle = [ |00\rangle + |10\rangle ] / \sqrt{2}$$

$$\mathbf{XOR}(\mathbf{H}|00\rangle) = [ |00\rangle + |11\rangle ] / \sqrt{2}$$

# Quantum Algorithms

# Quantum algorithms

Find the period of the function

$$f(x) = 1 + \cos(\pi x)$$

(Shor)

We work with two registers,  $X$  and  $Y$ :



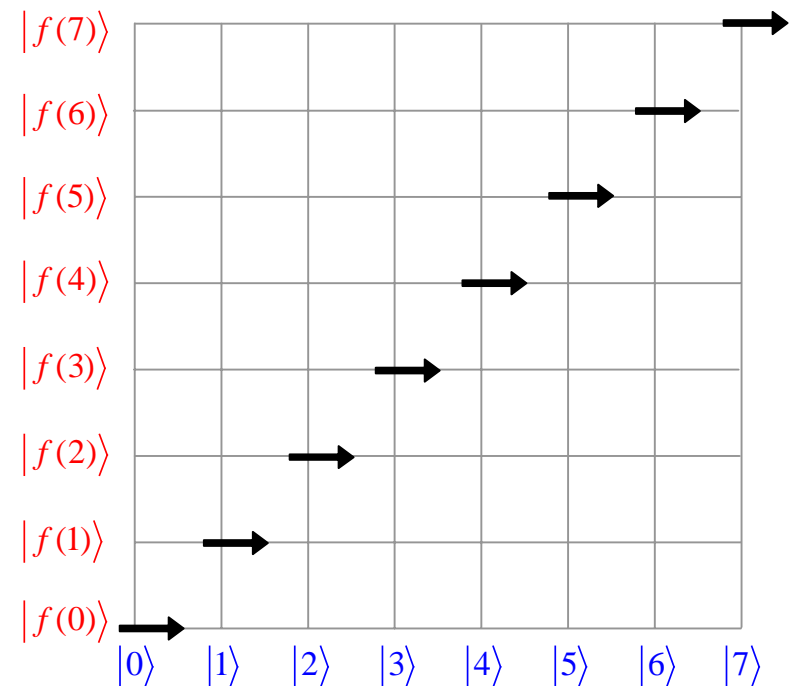
Place the  $X$  register into a superposition of all states,  $Y$  into the state  $0$ :

$$\begin{aligned} |\psi_X\rangle &= [ |000\rangle + |100\rangle + |010\rangle + |001\rangle + |011\rangle + |101\rangle + |110\rangle + |111\rangle ] / \sqrt{8} \\ &= [ |0\rangle + |1\rangle + |2\rangle + |3\rangle + |4\rangle + |5\rangle + |6\rangle + |7\rangle ] / \sqrt{8}, \end{aligned}$$

$$|\psi_Y\rangle = |000\rangle = |0\rangle.$$

Next, we entangle the  $X$  and  $Y$  registers, evaluating all function values  $f(x)$  in one go

$$\psi = \frac{1}{8} \sum_x |x, f(x)\rangle,$$



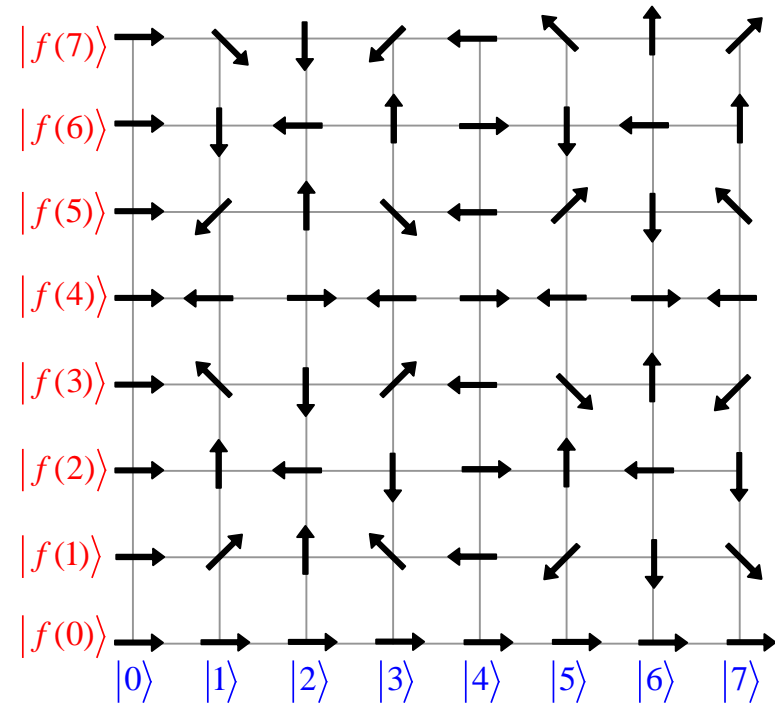
# Quantum algorithms, cont.

and then carry out a discrete Fourier transform on the  $X$  register,

$$|x\rangle \Rightarrow \frac{1}{\sqrt{8}} \sum_{k=0}^7 e^{2\pi i k x / 8} |k\rangle,$$

producing the superposition

$$\begin{aligned} |\psi\rangle &= \frac{1}{8} \sum_{x,k=0}^7 e^{2\pi i k x / 8} |k, f(x)\rangle \\ &= \frac{1}{8} |0\rangle [ |f(0)\rangle + |f(1)\rangle + \dots + |f(6)\rangle + |f(7)\rangle ] \\ &\quad + \frac{1}{8} |1\rangle [ |f(0)\rangle + e^{2\pi i / 8} |f(1)\rangle + \dots + e^{12\pi i / 8} |f(6)\rangle + e^{14\pi i / 8} |f(7)\rangle ] \\ &\quad + \frac{1}{8} |2\rangle [ |f(0)\rangle + e^{4\pi i / 8} |f(1)\rangle + \dots + e^{24\pi i / 8} |f(6)\rangle + e^{28\pi i / 8} |f(7)\rangle ] \\ &\quad + \dots \end{aligned}$$



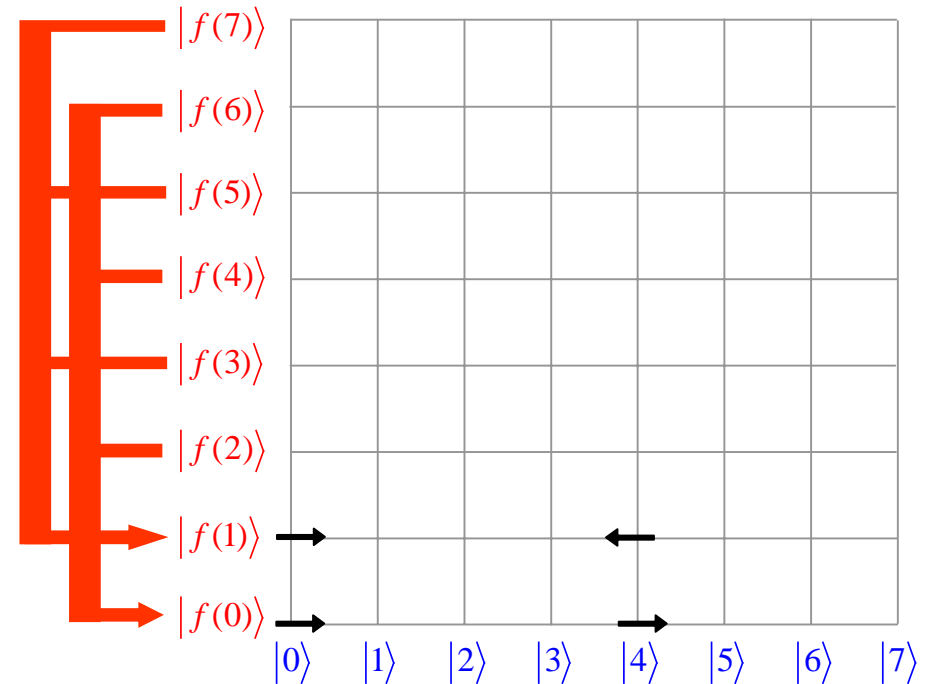
# Quantum algorithms, cont.

Finally, assume that  $f(x)$  has period 2,

$$f(0) = f(2) = f(4) = f(6),$$

$$f(1) = f(3) = f(5) = f(7);$$

then most amplitudes add up to **zero**:



$$\begin{aligned}
 |\psi\rangle = & \frac{1}{2}|0\rangle [ |f(0)\rangle + |f(1)\rangle ] \\
 & + \frac{1}{8}|1\rangle [ |f(0)\rangle [1 + e^{4\pi i/8} + e^{8\pi i/8} + e^{12\pi i/8}] + |f(1)\rangle [e^{2\pi i/8} + e^{6\pi i/8} + e^{10\pi i/8} + e^{14\pi i/8}] ] \\
 & + \frac{1}{8}|2\rangle [ |f(0)\rangle [1 + e^{8\pi i/8} + e^{16\pi i/8} + e^{24\pi i/8}] + |f(1)\rangle [e^{4\pi i/8} + e^{12\pi i/8} + e^{20\pi i/8} + e^{28\pi i/8}] ] \\
 & + \dots \\
 & + \frac{1}{8}|4\rangle [ |f(0)\rangle [1 + e^{16\pi i/8} + e^{32\pi i/8} + e^{48\pi i/8}] + |f(1)\rangle [e^{8\pi i/8} + e^{24\pi i/8} + e^{40\pi i/8} + e^{56\pi i/8}] ] \\
 & + \dots
 \end{aligned}$$

# Quantum algorithms, cont.

In the end we obtain the state

$$|\psi\rangle = \frac{1}{2} \left[ |0, f(0)\rangle + |0, f(1)\rangle + |4, f(0)\rangle - |4, f(1)\rangle \right]$$

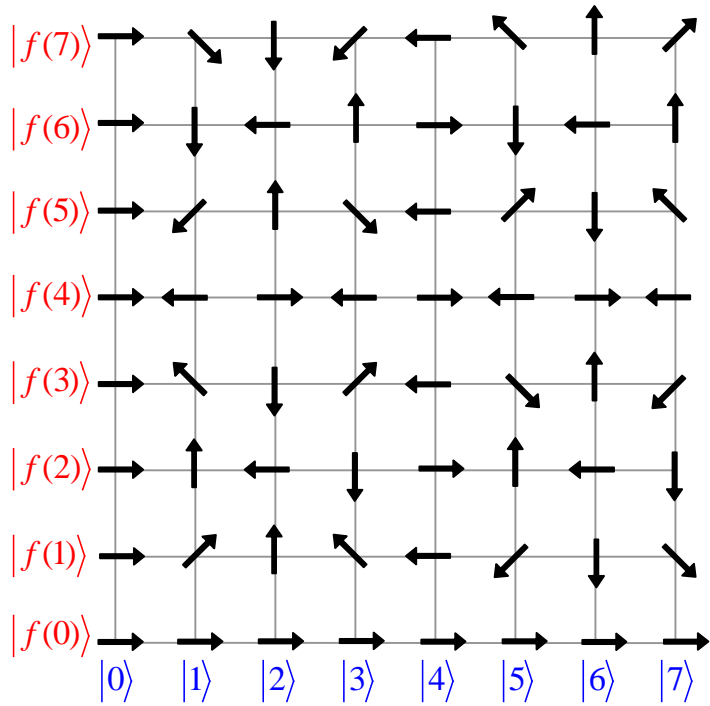
and a measurement of the  $X$  register always delivers a result  $k = 0$  or  $k = 4$ , each with probability  $1/2$ .

The period  $p$  of the function  $f$  then is given by

$$p = \text{maximal numerator} \left[ \text{red} \left( \frac{2^3}{k} \right) \right] = 2.$$

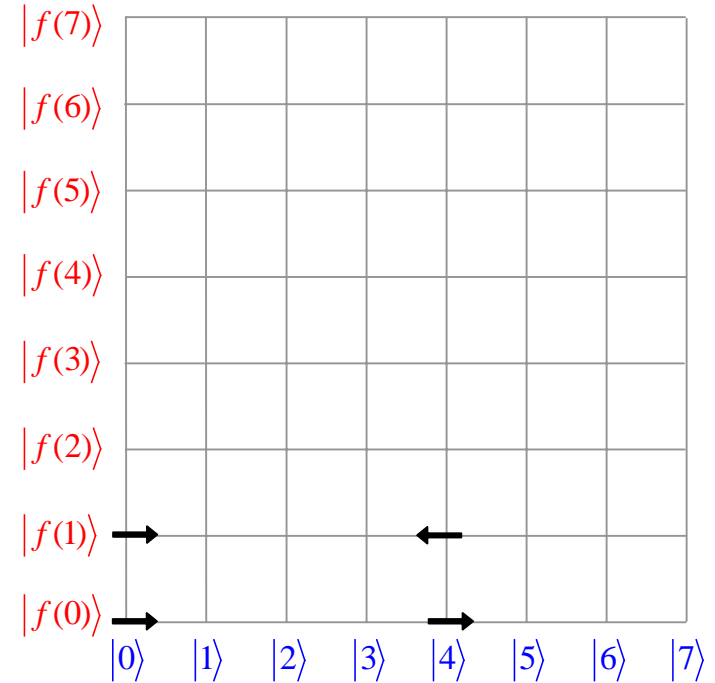
Once we know how to calculate the period of a function a few more steps are needed to factorize a number (Shor's algorithm, using **Euclid's algorithm** for finding the greatest common divisor of two numbers).

**Note, there is no quantum speed-up in adding numbers!  
Classical computers do a great job on that.**




'Cleanup'

due to interference



## Parallel evolution

Evaluate all function  $f(x)$  in one go



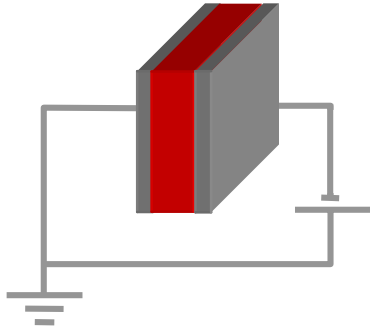
$$\psi = \frac{1}{8} \sum_x |x, f(x)\rangle$$



Hardware

# Hardware

Classical computer



Capacitors:

Bits

0:  $V = 0$ ,

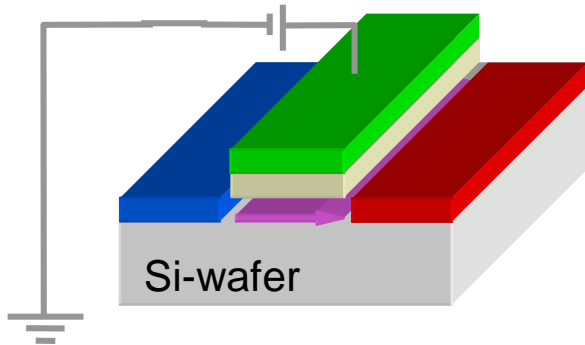
1:  $V > 0$ .

Transistors:

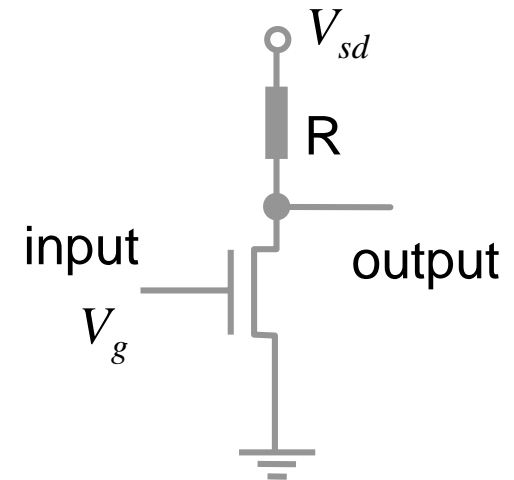
Gates

$V_g = 0$ , closed,

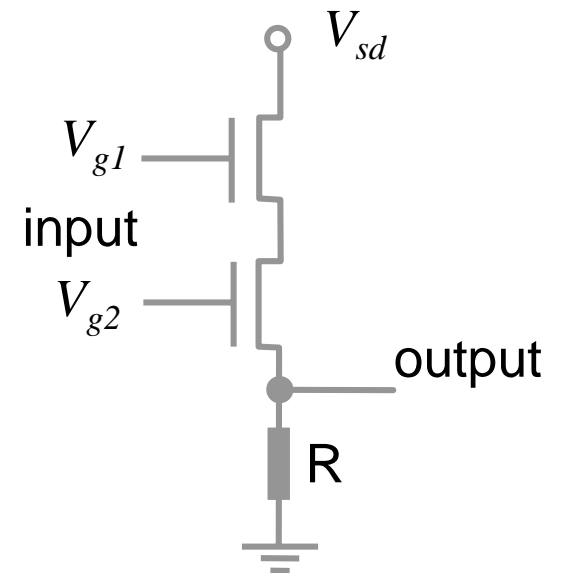
$V_g > 0$ , open.



1-bit gate: NOT



2-bit gate: AND



# Network model of quantum computing

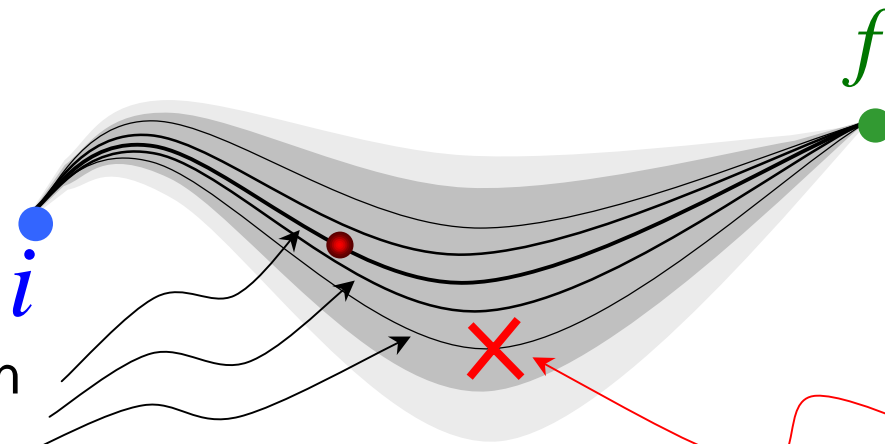
(David Deutsch, 1985)

initial state



- each qubit can be **prepared** in some known state,  $|00000\dots0000\rangle$ .
- each qubit can be **measured** in a basis,  $|01100\dots1010\rangle$ .
- the qubits can be **manipulated** through quantum gates
- the qubits are **protected** from decoherence

final state



Parallel evolution  
providing the  
quantum speedup.



Perturbations  
from the environment  
destroy the parallel evolution of the computation

# Physical implementations

## Quantum optics, NMR-schemes

Good decoupling & precision:

- trapped atoms (Cirac & Zoller)
- photons in QED cavities (Monroe ea, Turchette ea)
- molecular NMR (Gershenfeld & Chuang)
- $^{31}\text{P}$  in silicon (Kane)

## Solid state implementations

Good scalability & variability:

- spins on quantum dots (Loss & DiVincenzo)
- $^{31}\text{P}$  in silicon (Kane)
- Josephson junctions, charge (Schön ea, Averin)  
phase (Bocko ea, Mooij ea)



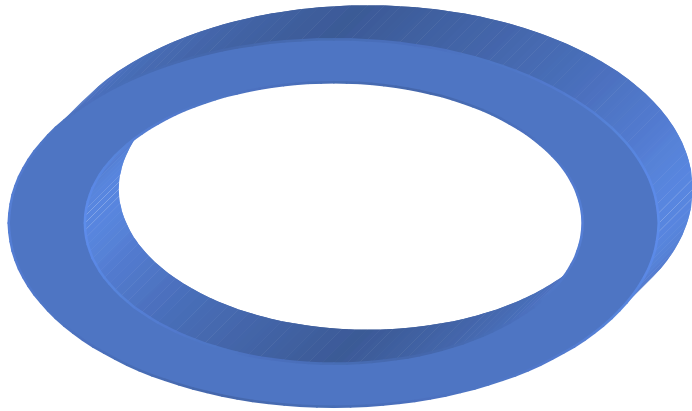
All hardware implementations of quantum computers have to deal with the conflicting requirements of **controllability** while minimizing the coupling to the environment in order to **avoid decoherence**.



Have to deal with individual atoms, photons, spins,.....  
Problems with control, interconnections, measurements.

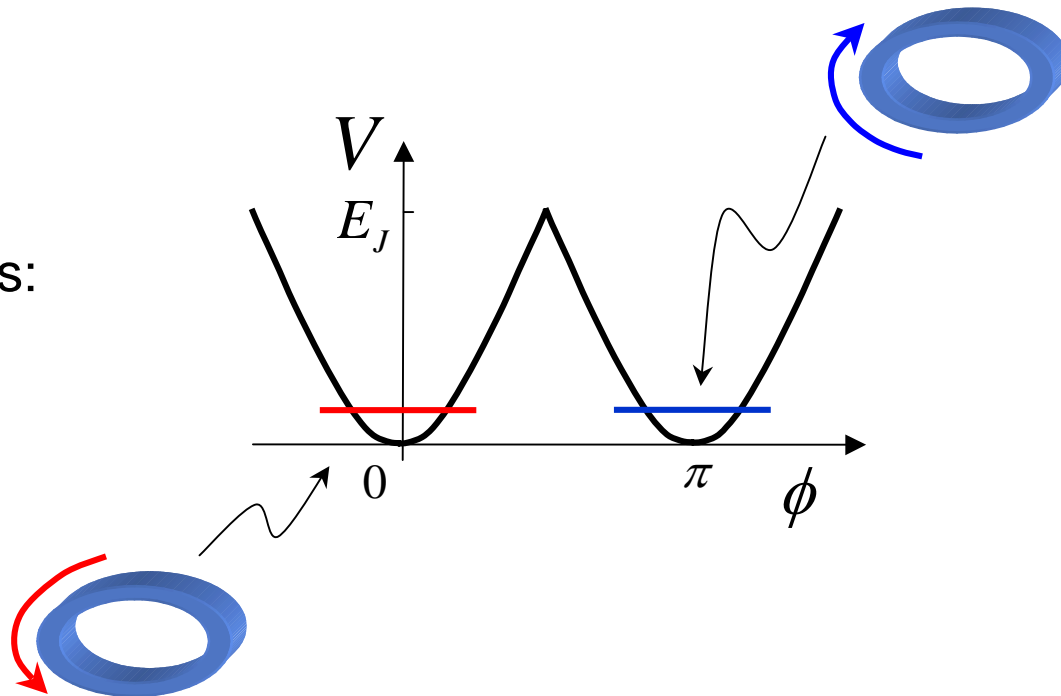
Have to deal with many degrees of freedom.  
Problems with decoherence.

# Superconducting Phase Qubits

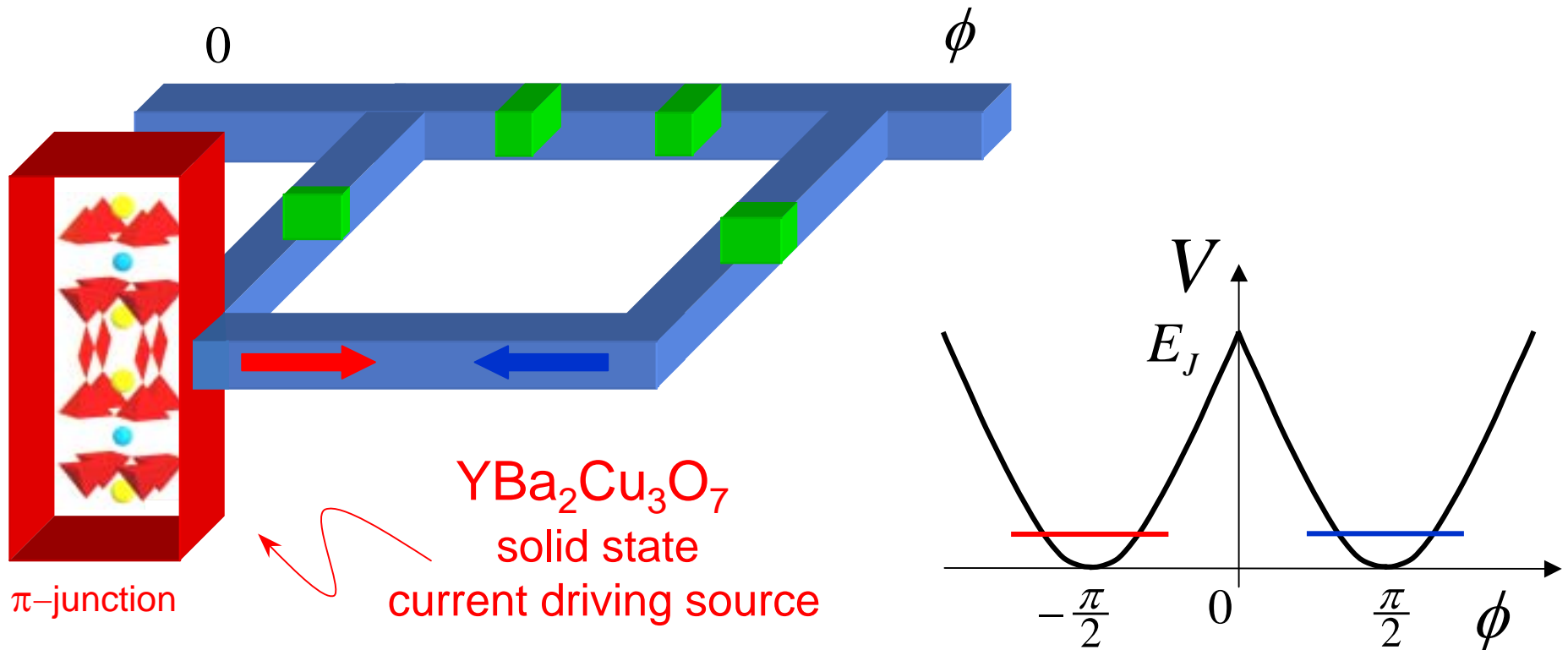


In a superconducting ring, currents do not decay: persistent current states.

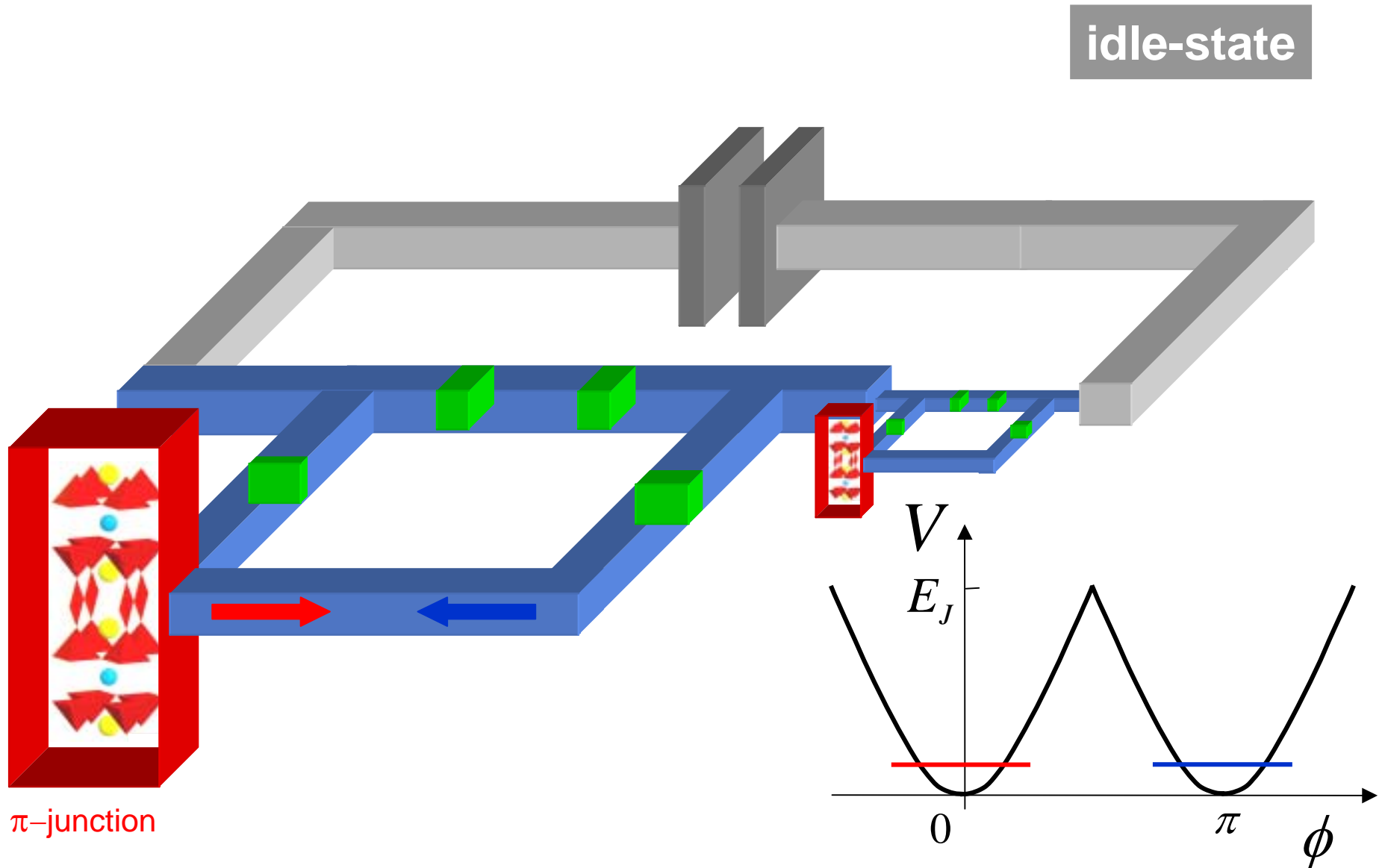
Construct a qubit with 2 persistent current states:



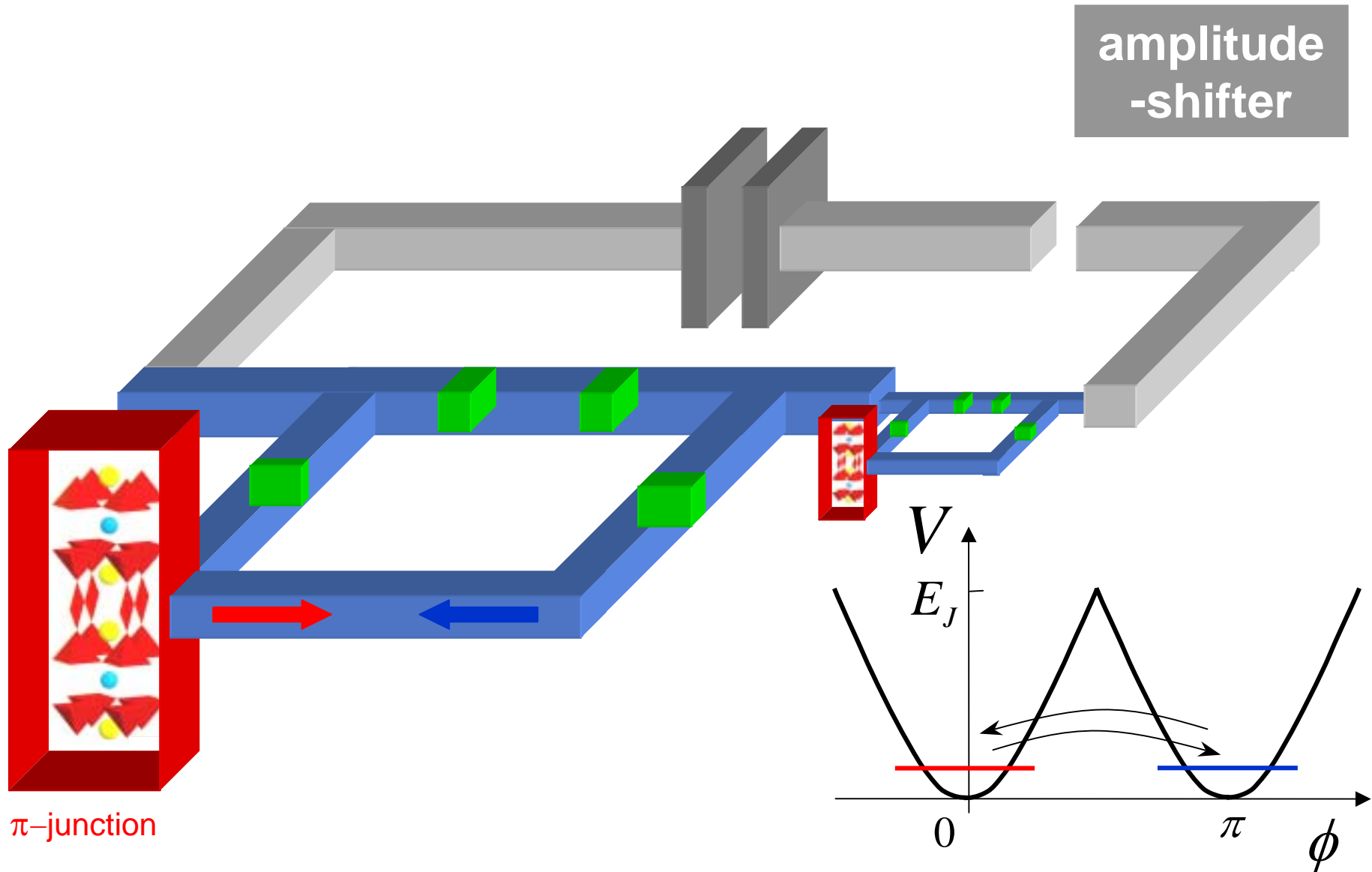
# Superconducting Phase Qubits



# Superconducting Phase Qubits

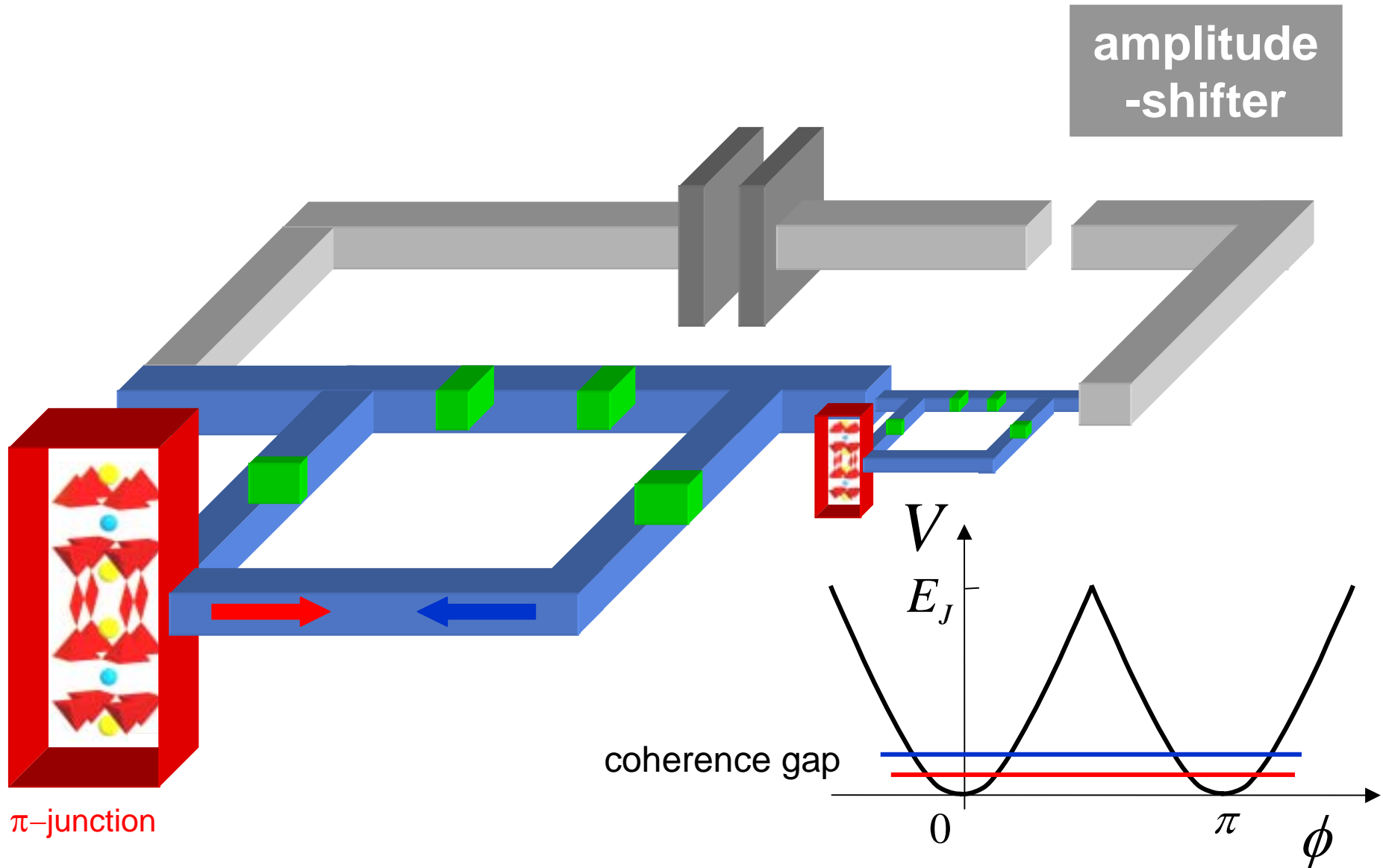


# Superconducting Phase Qubits

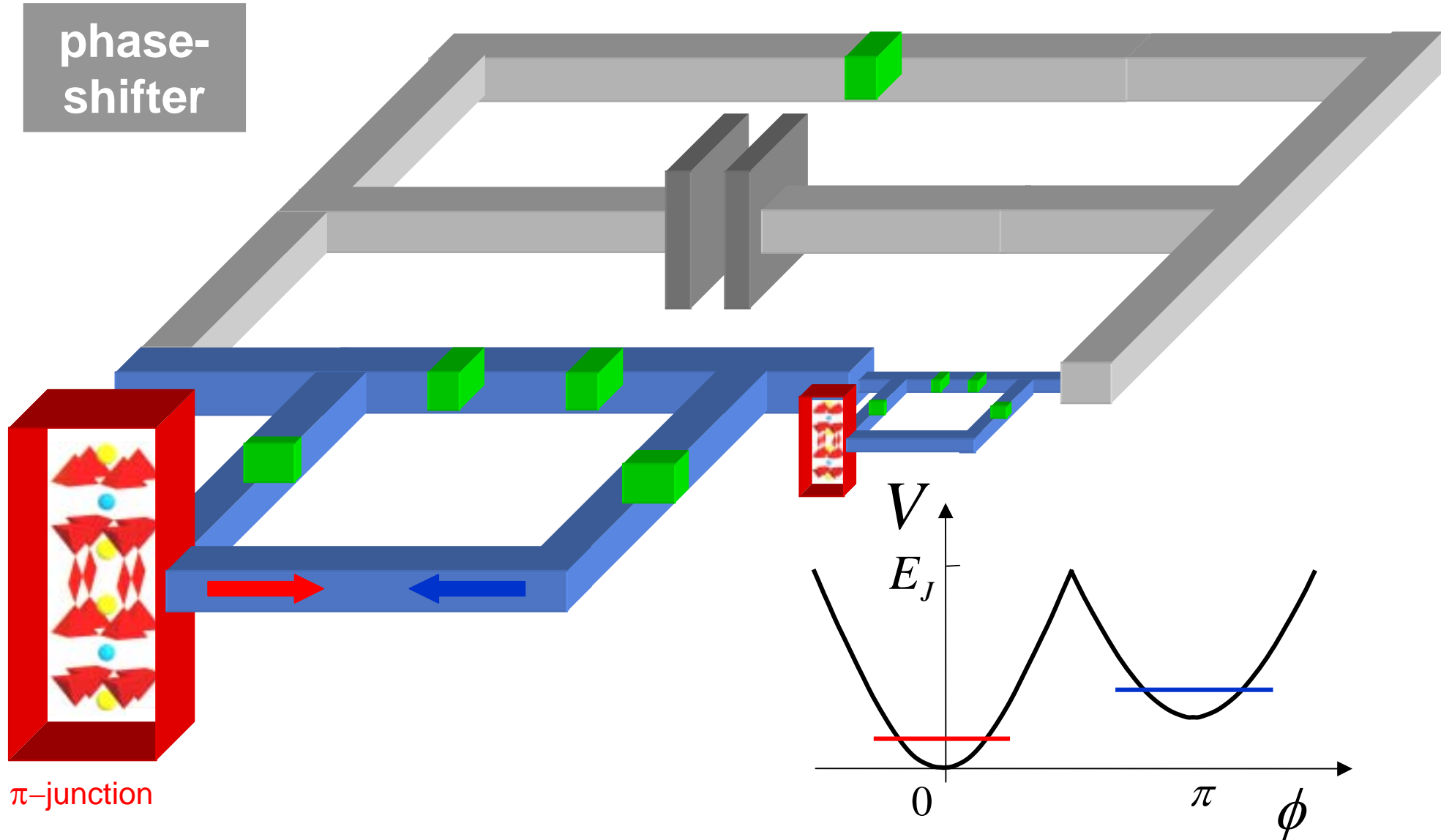




# Superconducting Phase Qubits



# Superconducting Phase Qubits



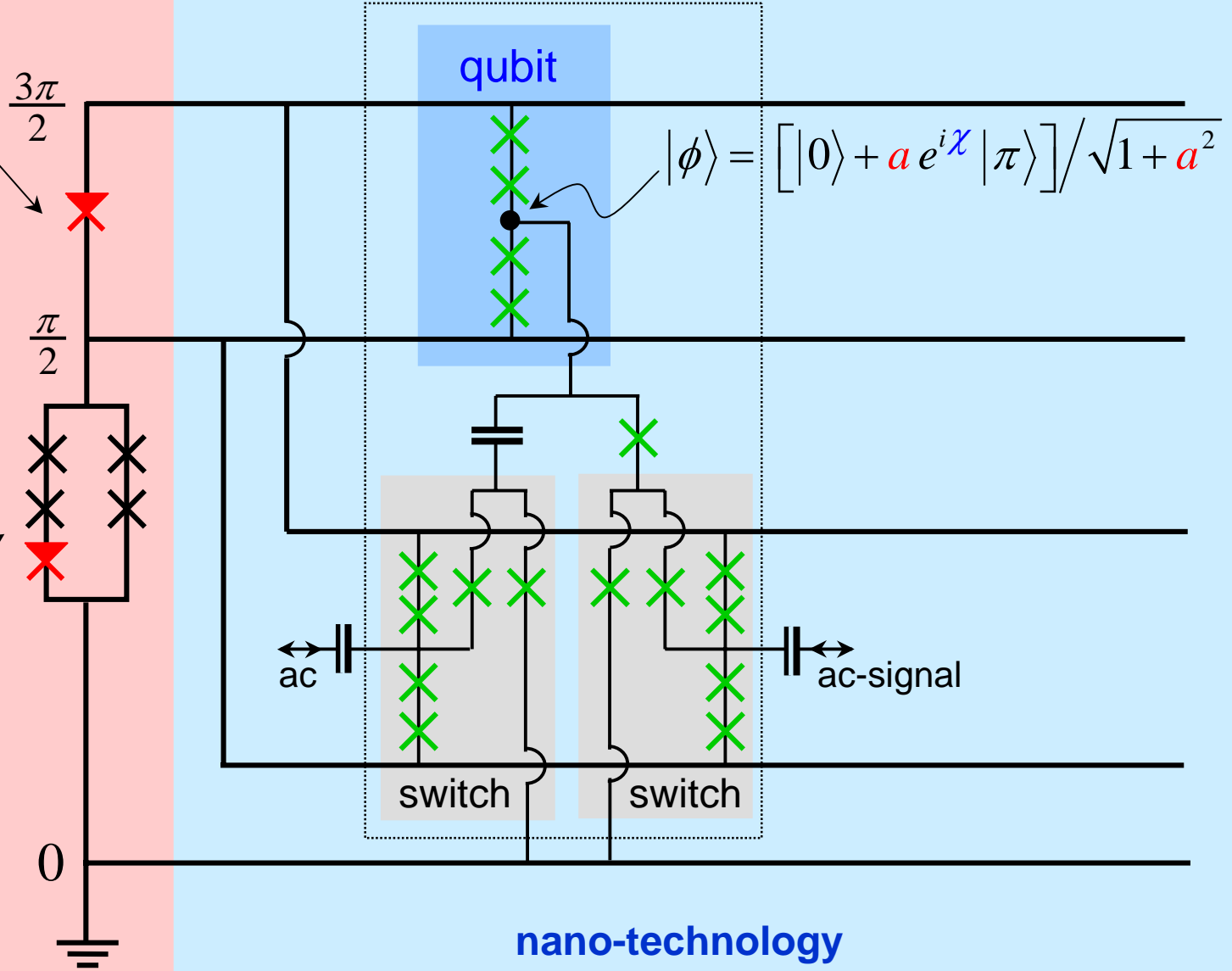
# 'Blueprint'



$\pi$ -phase shifter  
(~battery)

$\pi/2$ -phase shifter  
(~voltage divider)

**materials science**  
provides the  
driving sources  
for the qubit



**nano-technology**  
provides the  
quantum nature of the qubit

# Quantum Computation

## Classical computer:

- the information is stored in classical bits, values 0,1
- usual operations NOT, AND, OR
- general purpose device



## Quantum computer:

- the information is stored in quantum two-level systems, **qubits**
- **unitary operations**, single- and two-qubit operations (XOR)
- powerful in calculating **specific tasks**

### Algorithms:

- prime factorization (Shor)
- data base search (Grover)
- random number generator
  - quantum simulation

### Applications:

- quantum teleportation
- quantum cryptography

**Quantum computation** is an interdisciplinary field, with contributions from mathematics, (theoretical) computer science, physics, and chemistry

### Implementations:

- Quantum optics
- **Solid state**
- NMR

- **History:** from mechanics to nanoelectronics
- **Information Theory:** Turing machines & complexity
- **Quantum Mechanics:** superpositions & entanglement
- **Quantum Games:** no-cloning, cryptography, teleportation
- **Quantum Bits and Gates**
- **Quantum Algorithms:** Shor's period finder
- **Hardware:** superconducting phase qubits



## Thanks to

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Mikhail Feigel'man

Oskar Hallatschek

Akira Tonomura

Ulrich Helg

Hanni Hediger

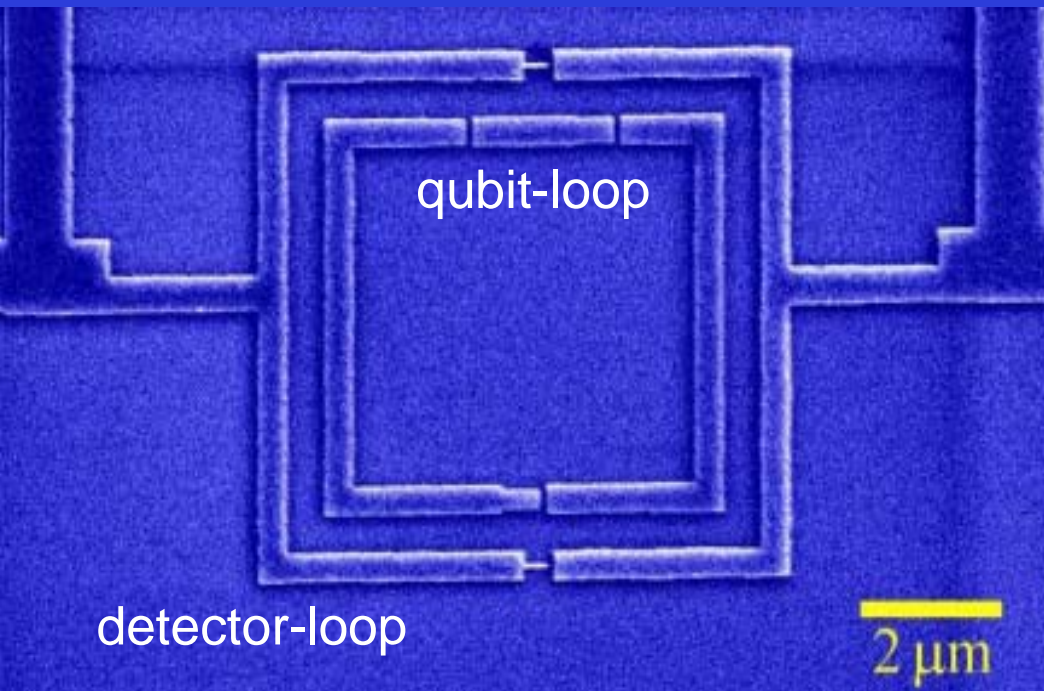
Elmar Heeb



# Solid State Qubits: Recent Achievements

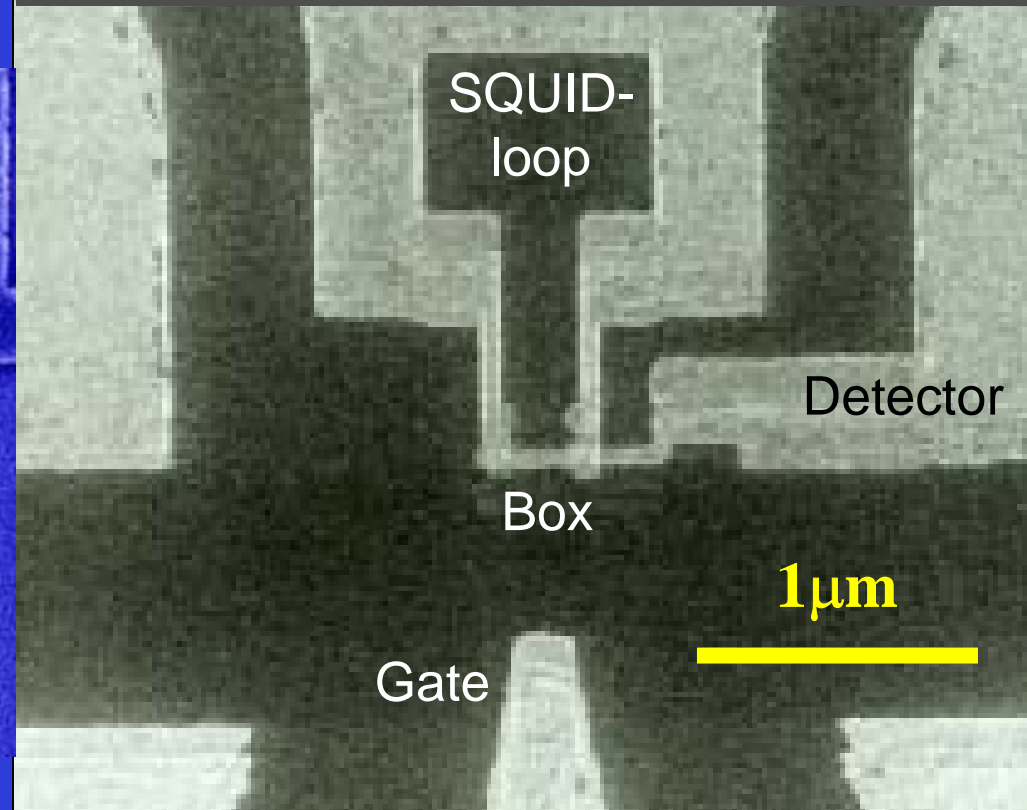
## Superconducting Phase Qubit

Al-technology,  
observation of coherence gap  
via *rf*-absorption;  
Hans Mooij *et al.*, TU-Delft



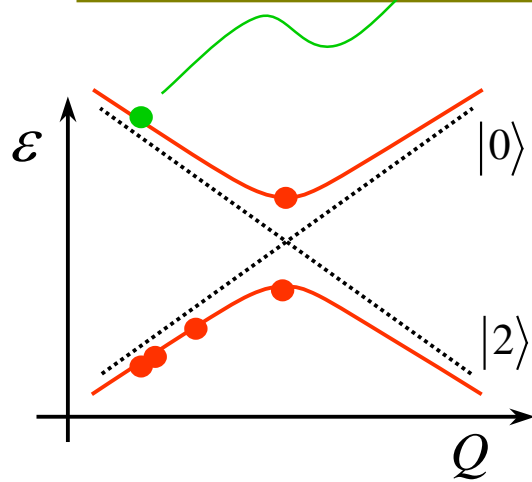
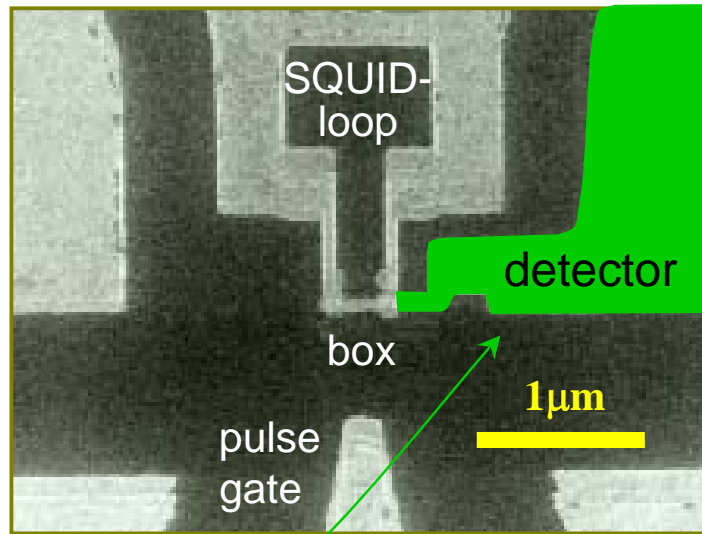
## Superconducting Charge Qubit

Al-technology,  
observation of coherence gap  
via *rf*-absorption and via a  
real time experiment;  
J. Tsai *et al.*, NEC Tsukuba



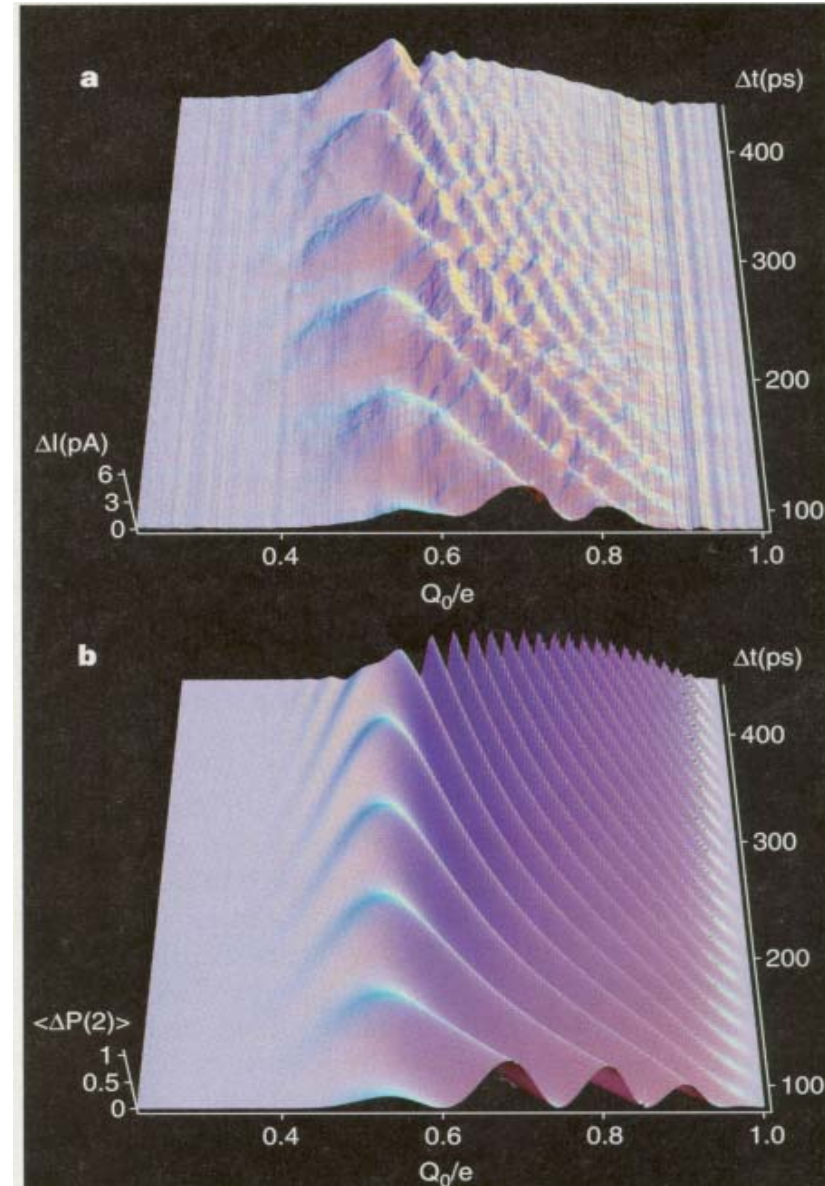


# Charge (Nakamura *et al.*, 1999)



This time domain experiment shows coherent charge oscillations of **50 – 100 ps** duration during a total coherence time of **2 ns**.

# Coherent devices



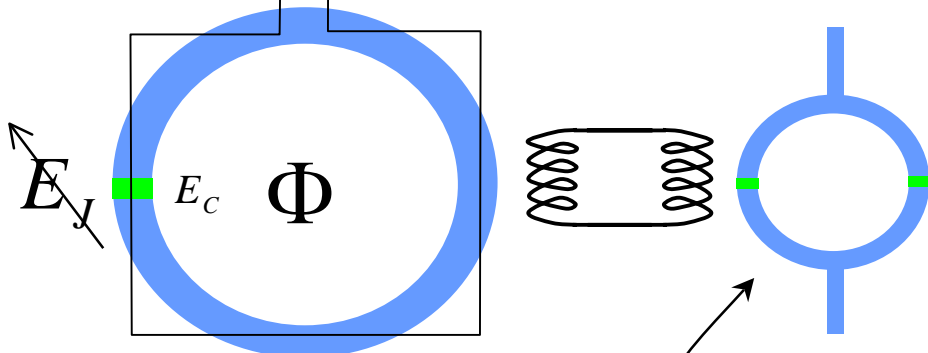


Flux (Friedman *et al.*, 2000)

# Coherent devices

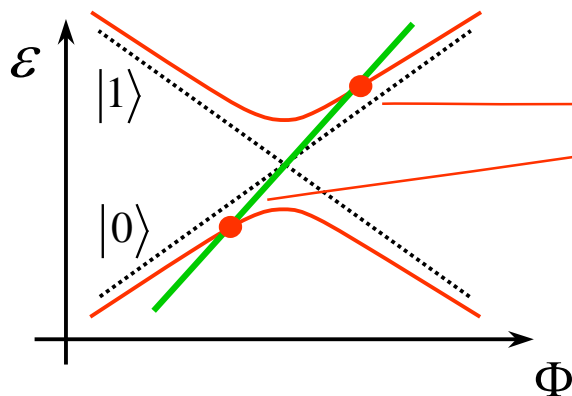
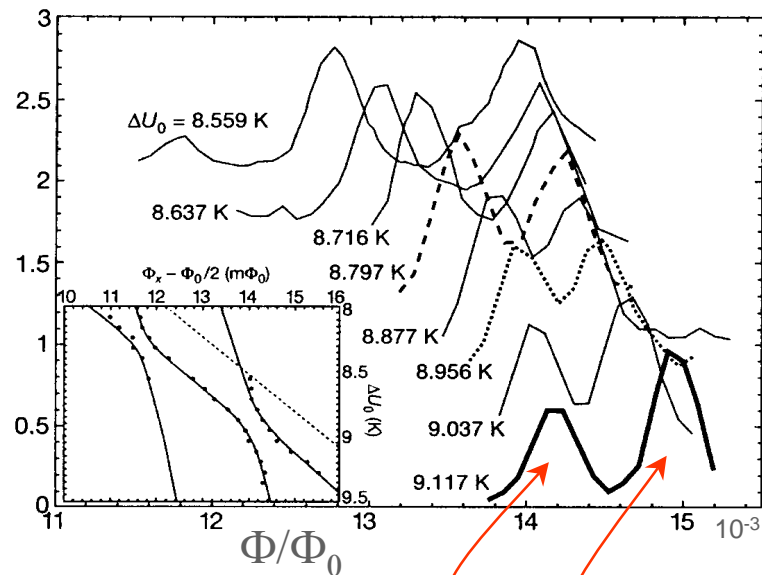
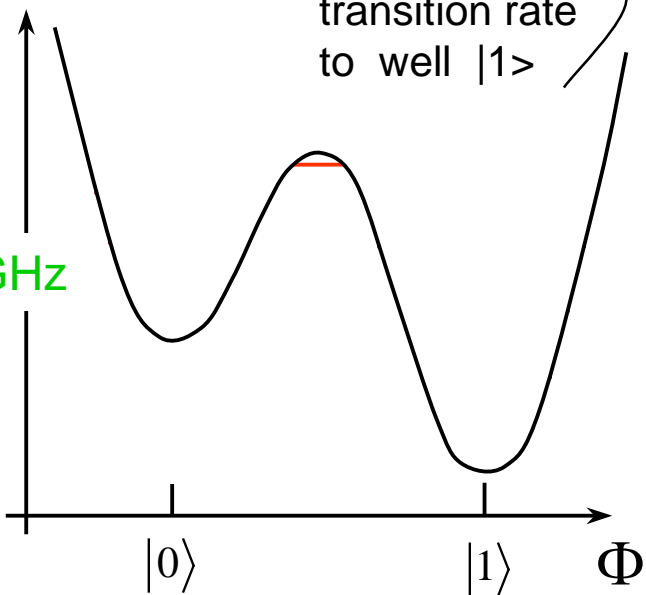
rf - SQUID

$I_g$

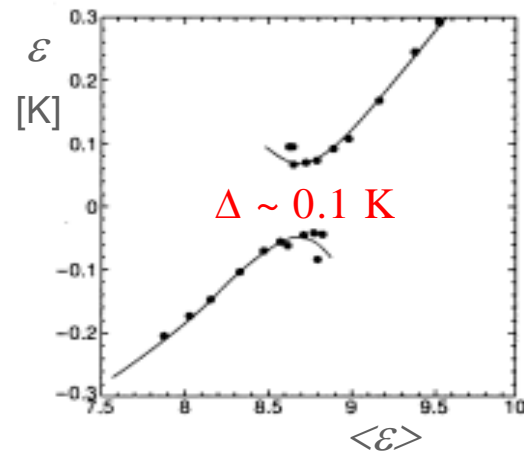


detect transition rate to well  $|1\rangle$

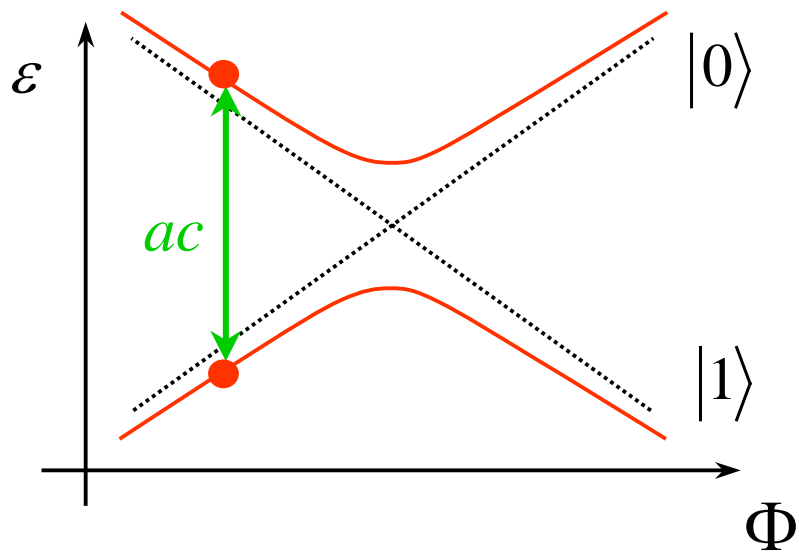
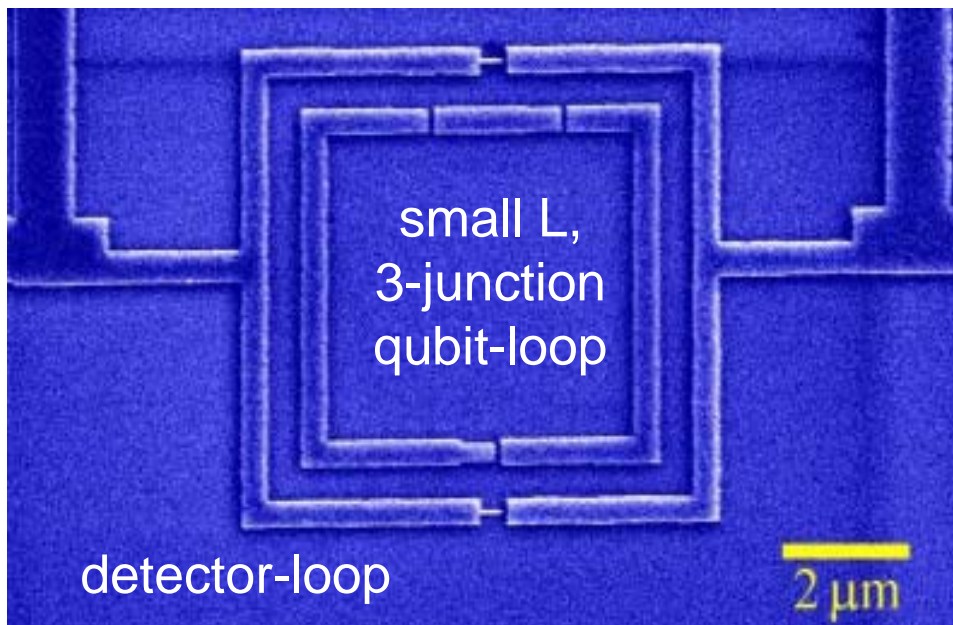
96 GHz



This experiment demonstrates the coherent current oscillation of  $\sim 10^9$  Cooper pairs producing an oscillating moment of size  $\sim 10^{10} \mu_B$ .

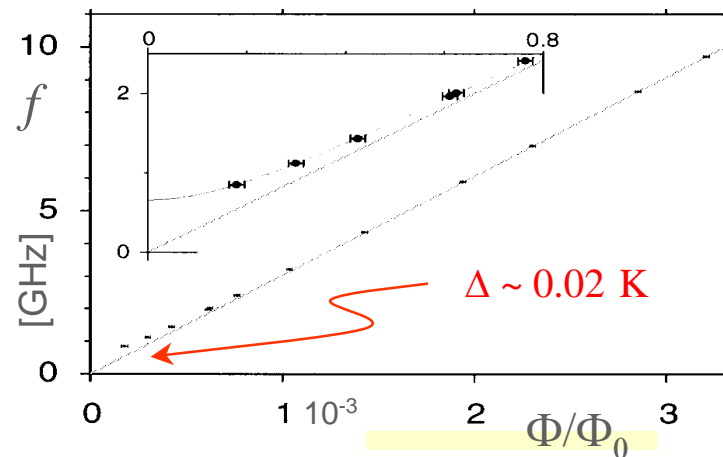
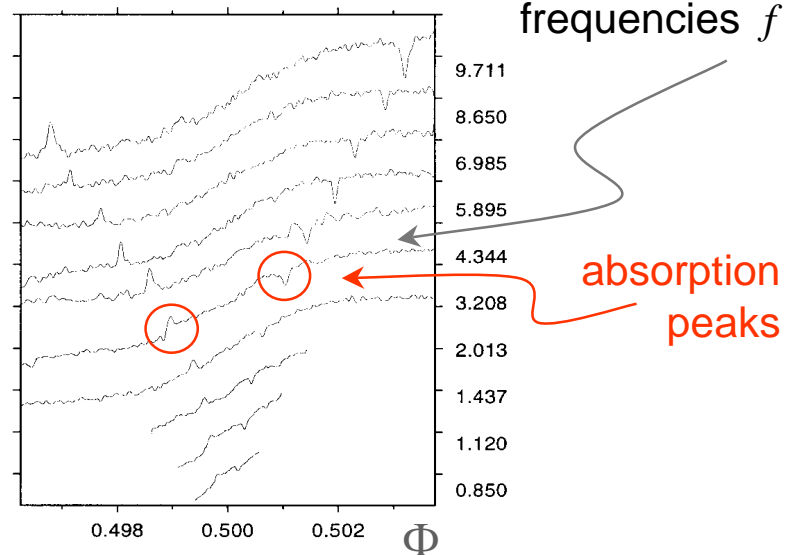


Phase (Mooij *et al.*, 2000)



## Coherent devices

Scanning flux  $\Phi$  at different microwave frequencies  $f$



This experiment measures coherent current (0.5  $\mu\text{A}$ ) oscillations over  $\sim 5$  cycles.

